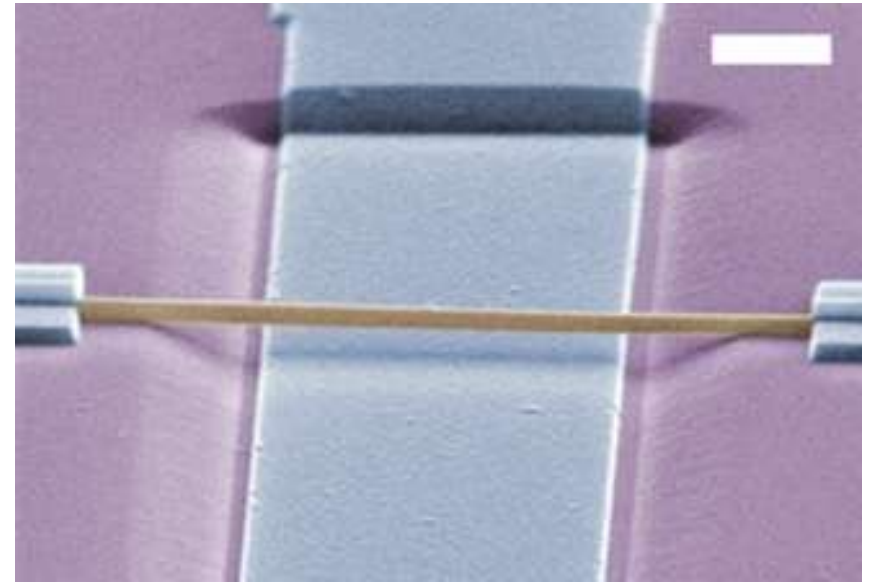
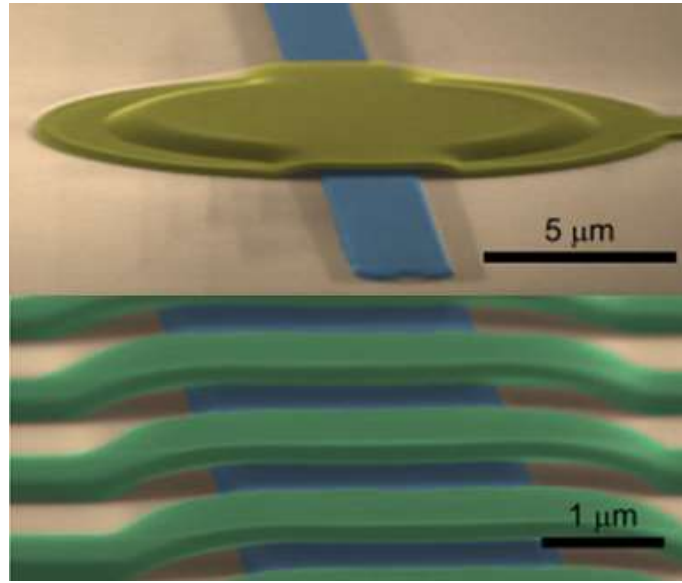
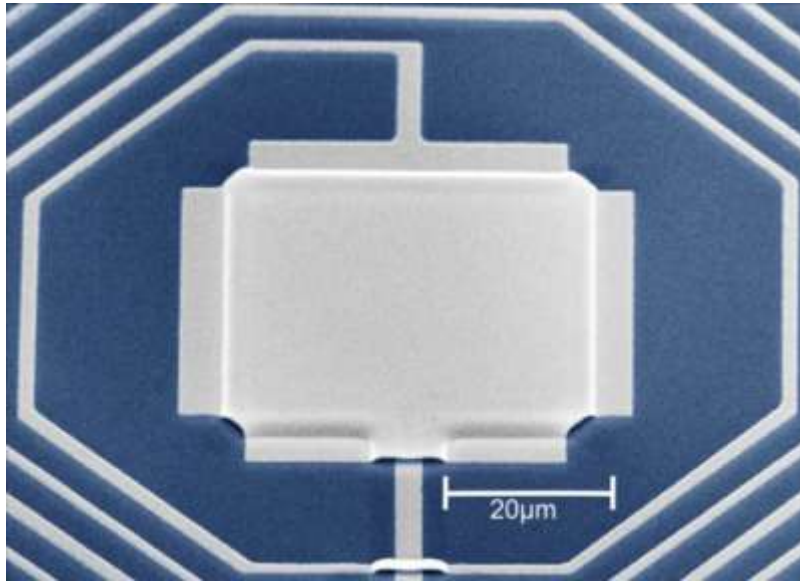




중시계 진동자의 광역학적 상호작용

Optomechanical interactions in mesoscopic oscillators

The 14th School of Mesoscopic Physics: Mesoscopic Interactions

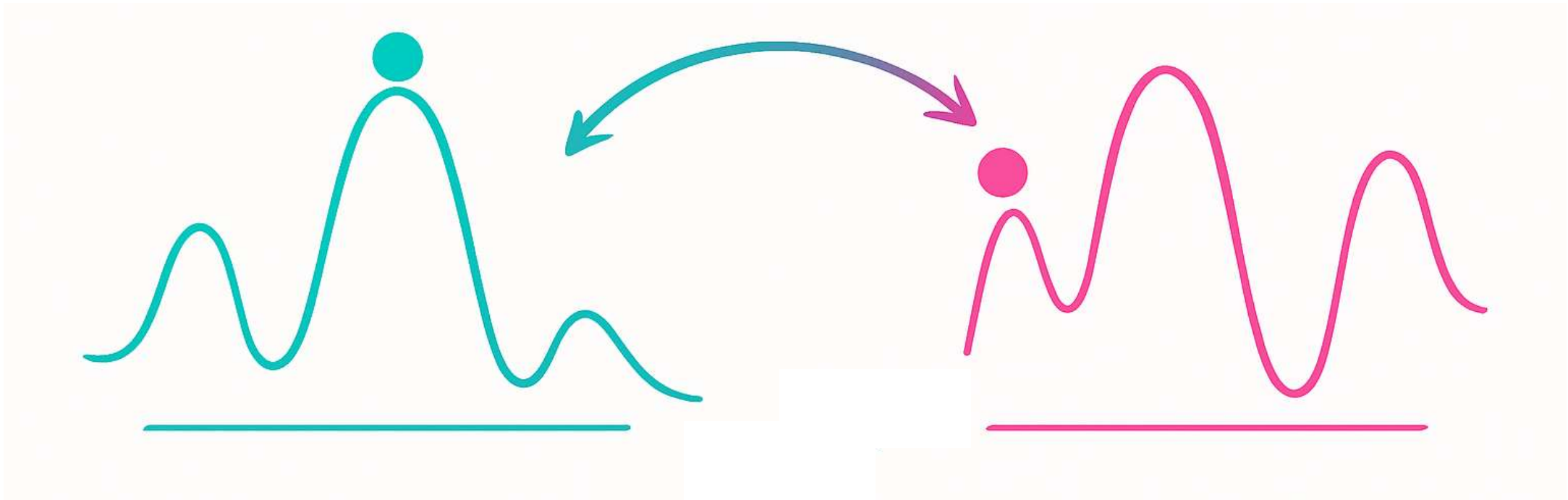
TOPICS

1. Quantum Transport and Topological Matters
2. Light-Matter Interaction ✓
3. Mesoscopic Superconductor interaction
4. Phonon-Cooper pair interaction
5. Josephson Diode

ORGANIZERS

Myoung-Ho Bae (KRISS)
Hyungkook Choi (Jeonbuk Nat'l Univ.)
Hyoungsoon Choi (KAIST)
Yong-Joo Doh (GIST)
Nojoon Myoung (Chosun Univ.)
Myunglae Jo (KNU)
Minkyung Jung (DGIST)
Hee Chul Park (Pukyong Nat'l Univ.)
Heungsun Sim (KAIST)
Sekwon Kim (KAIST)

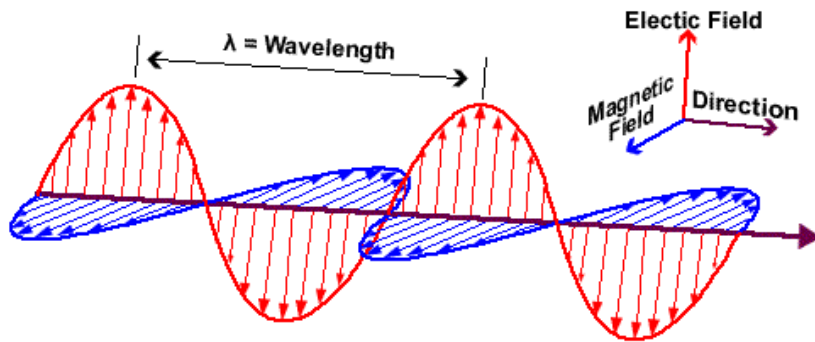
상호(相互) 작용 - (Mutual) Interaction



* “interacting quantum states” image generated by chatGPT

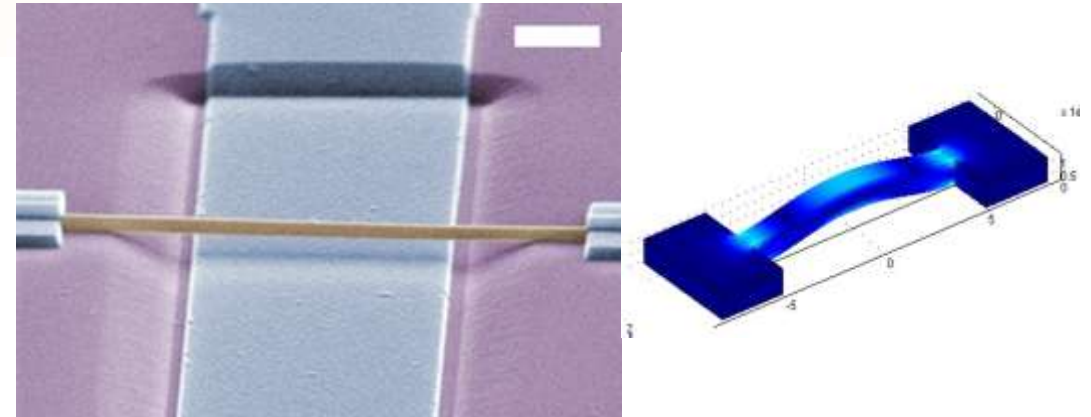
광역학적 상호작용

빛



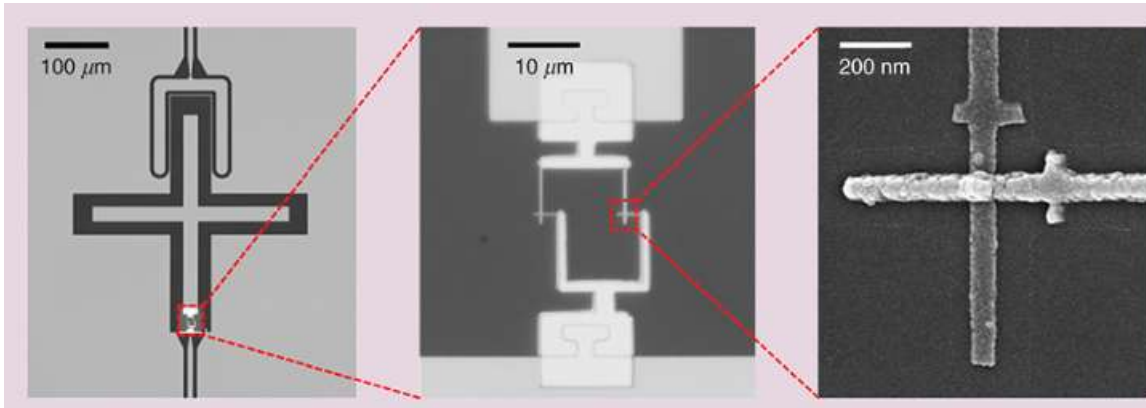
마이크로파 대역 전자기파
(0.3~30 GHz)

중시계 (역학적) 진동자

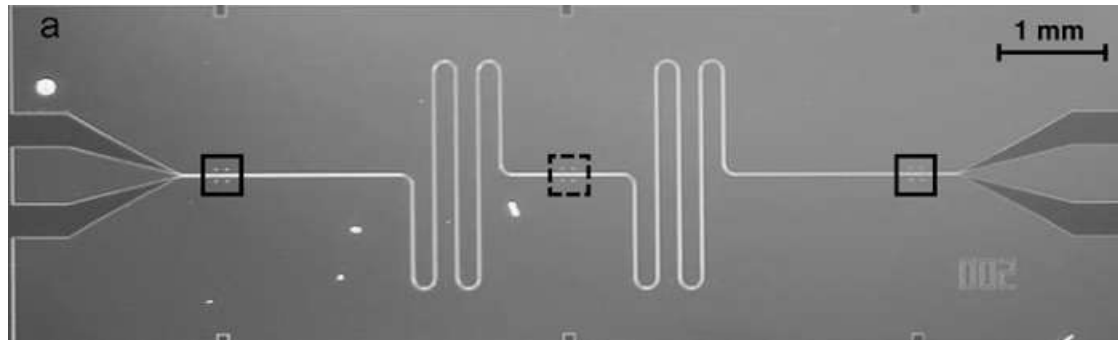


마이크로/나노 역학적 진동자
(MHz~GHz)

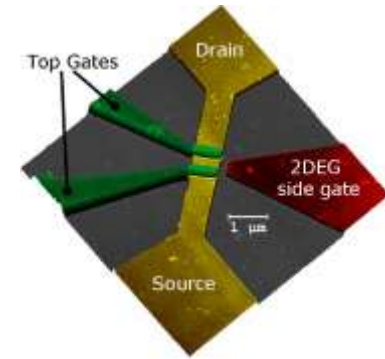
마이크로파 중시계 소자의 예



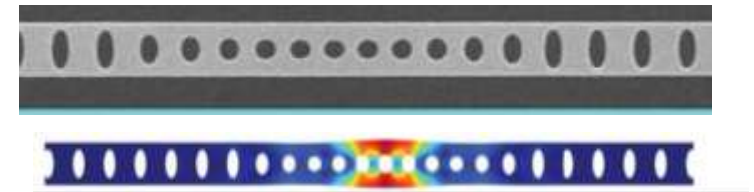
Superconducting qubit



Superconducting CPW resonator



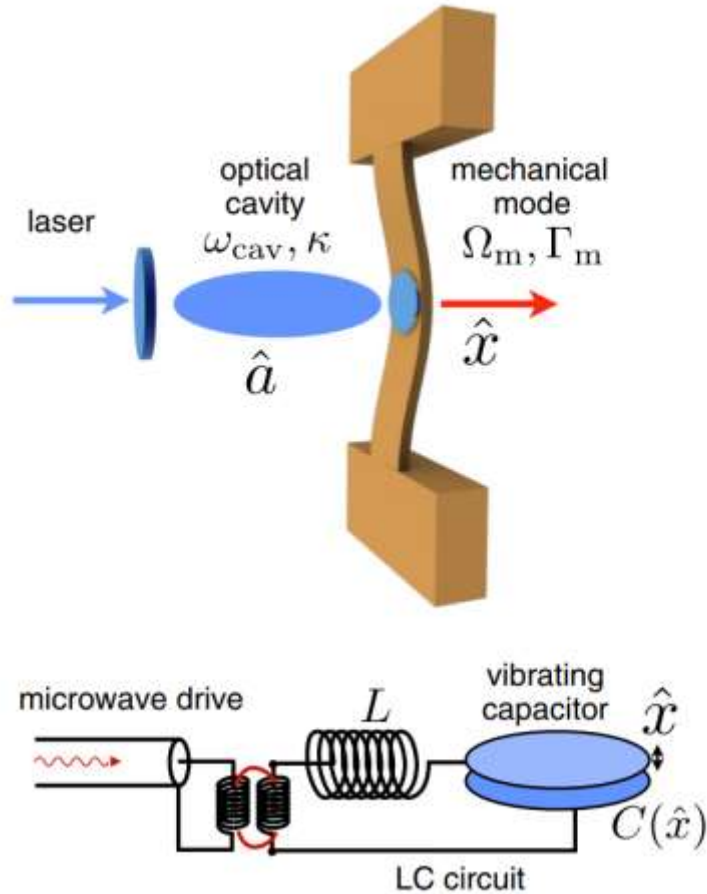
Electron quantum dot qubit



Nano-acoustic resonator

*** Low temperature necessary:
Why? Microwave photon energy > thermal energy
e.g.) 1 GHz microwave photon ~ 50 mK thermal energy

Cavity optomechanical system (광역학계)



[Cavity Hamiltonian]

$$= \hbar \omega_{cav} \hat{a}^\dagger \hat{a}$$

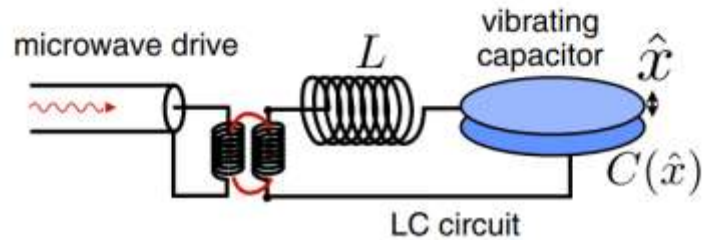
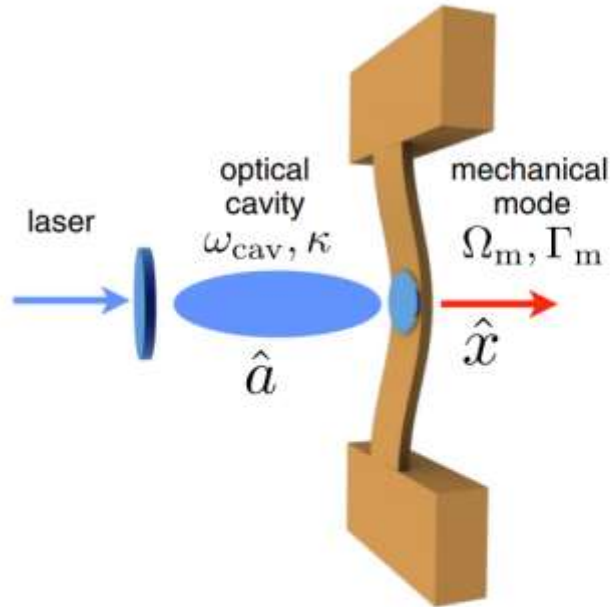


[Optomechanical interaction Hamiltonian]

$$= \hbar \frac{\partial \omega_{cav}}{\partial x} \hat{x} \hat{a}^\dagger \hat{a} \quad \text{“radiation pressure”}$$

* Aspelmeyer *et al.*, *Rev. Mod. Phys.* **86**, 29 (2014).

Cavity optomechanical system (광역학계)



[Cavity Hamiltonian]

$$= \hbar \omega_{cav} \hat{a}^\dagger \hat{a}$$



[Optomechanical interaction Hamiltonian]

$$= \hbar \frac{\partial \omega_{cav}}{\partial x} \hat{x} \hat{a}^\dagger \hat{a} \quad \text{"radiation pressure"}$$

$$\Rightarrow L_{cav} = n\lambda = n \frac{2\pi c}{\omega_{cav}} \rightarrow \delta x = \delta L_{cav} \approx -n \frac{2\pi c}{\omega_{cav}^2} \delta \omega_{cav}$$

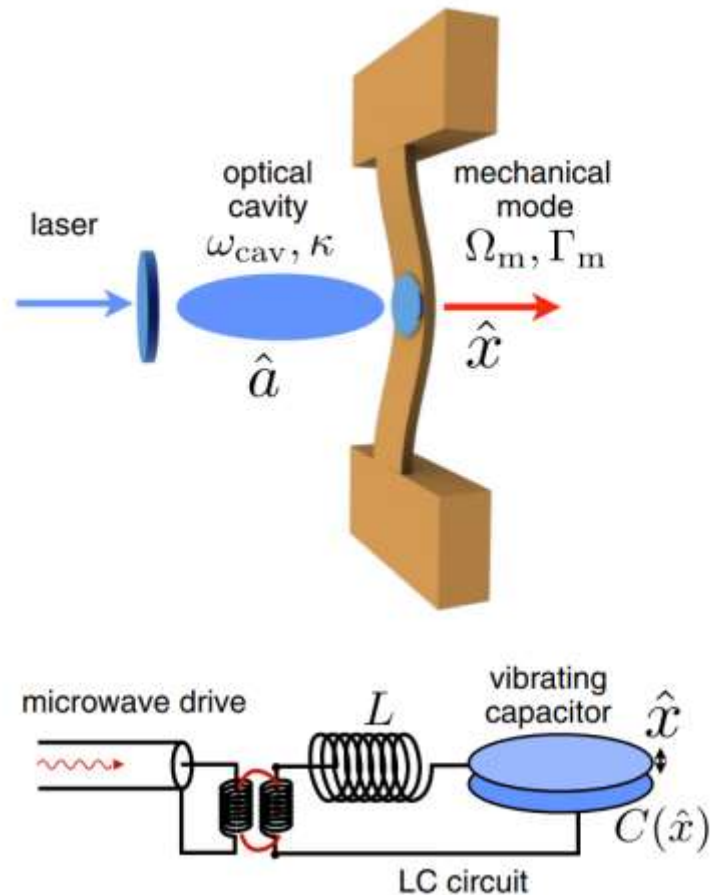
$$\frac{\partial \omega_{cav}}{\partial x} \approx -\frac{\omega_{cav}^2}{2\pi c n}$$

$$\Rightarrow \omega_{cav} = \frac{1}{\sqrt{LC}}; C = \frac{\epsilon_0 A_{cap}}{d_{cap}} \rightarrow \delta x = \delta d_{cap} \approx \frac{2d_{cap}}{\omega_{cav}} \delta \omega_{cav}$$

$$\frac{\partial \omega_{cav}}{\partial x} \approx \frac{\omega_{cav}}{2d_{cap}}$$

* Aspelmeyer et al., Rev. Mod. Phys. **86**, 29 (2014); Devoret et.al., lecture notes of les houches summer school (2011).

Cavity optomechanical system (광역학계)



Mechanical oscillator couples to photons

$$\hat{H} = \hbar\omega_{cav}\hat{a}^\dagger\hat{a} + \hbar\Omega\hat{b}^\dagger\hat{b} + \hbar g_0\hat{a}^\dagger\hat{a}(\hat{b}^\dagger + \hat{b})$$

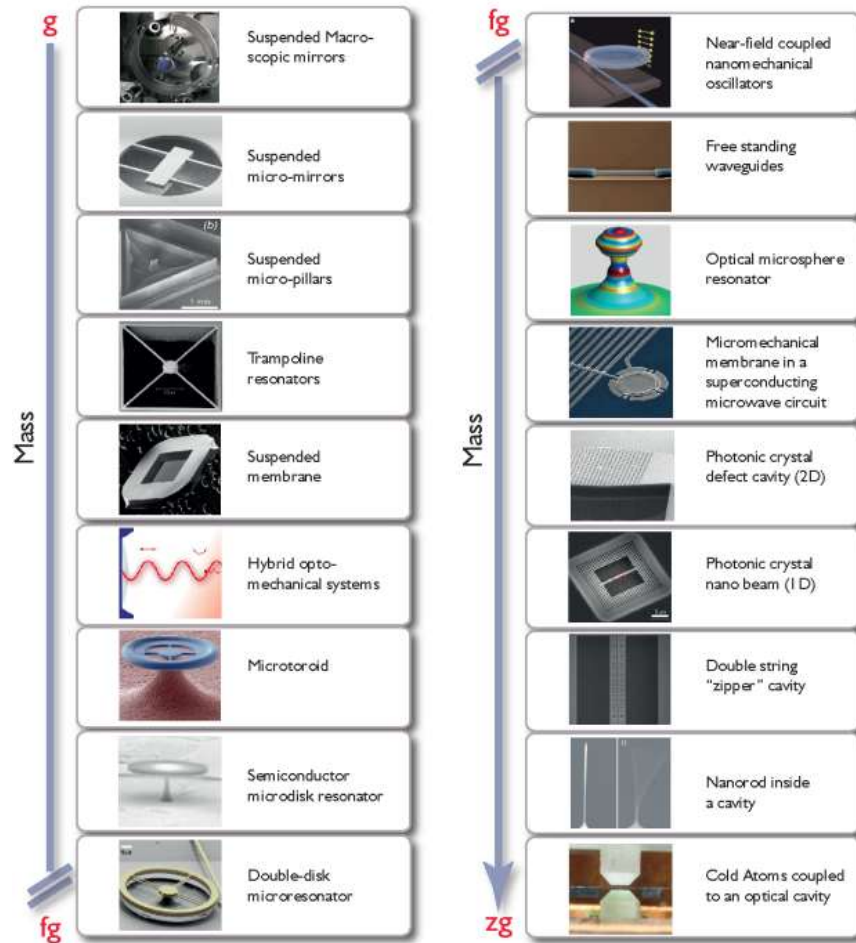
↑ photon
 ↑ phonon
 ↑ interaction

* Optomechanical “single-photon” coupling strength

$$g_0 = \frac{\partial \omega_{cav}}{\partial x} x_{zpf}$$

* Aspelmeyer *et al.*, *Rev. Mod. Phys.* **86**, 29 (2014); Devoret *et.al.*, *lecture notes of les houches summer school* (2011).

Cavity optomechanical system (광역학계)



Mechanical oscillator couples to photons

$$\hat{H} = \hbar\omega_c \hat{a}^\dagger \hat{a} + \hbar\omega_m \hat{b}^\dagger \hat{b} + \hbar g \hat{a}^\dagger \hat{a} (\hat{b}^\dagger + \hat{b})$$

↑ photon
 ↑ phonon
 ↑ interaction

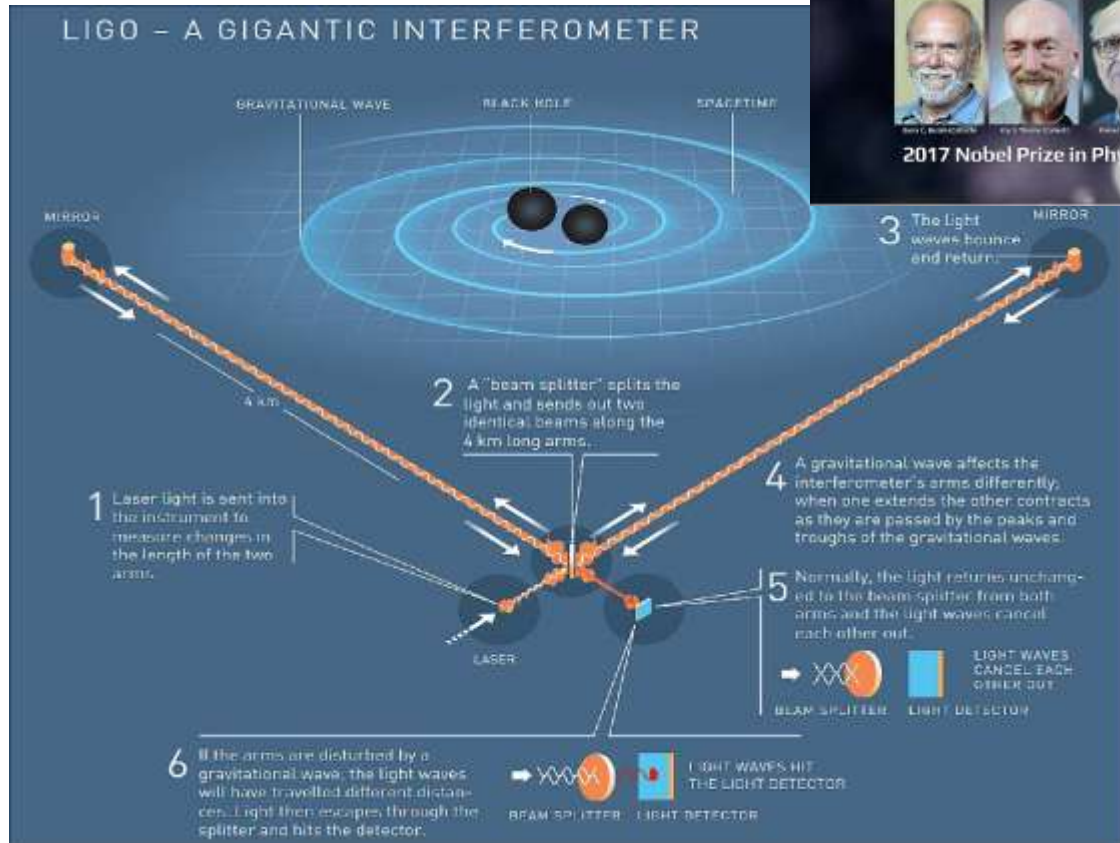
$= \hat{x}/x_{zp}$

- Ultrasensitive measurements
- Quantum hybrid systems
- Fundamental tests of quantum mechanics with gravity
- Classical/Quantum information processing ...

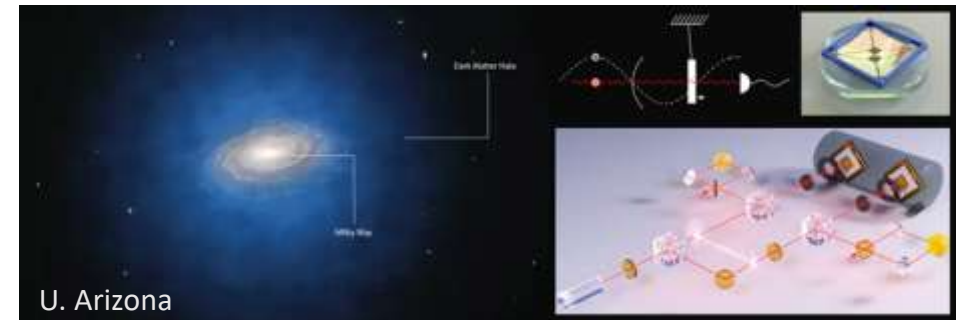
* Aspelmeyer *et al.*, *Rev. Mod. Phys.* **86**, 29 (2014).

Cavity optomechanical system (광역학계)

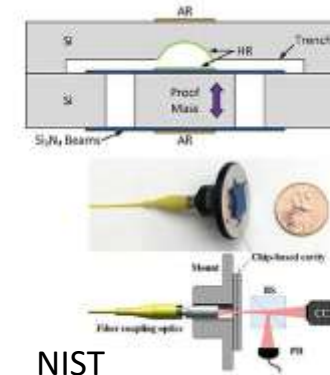
중력파



암흑물질

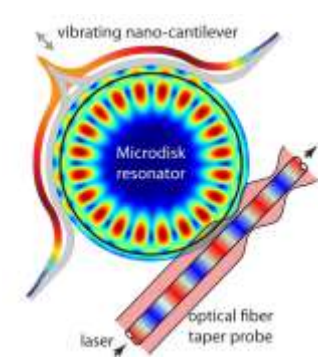


가속도



NIST

Atomic force



NIST

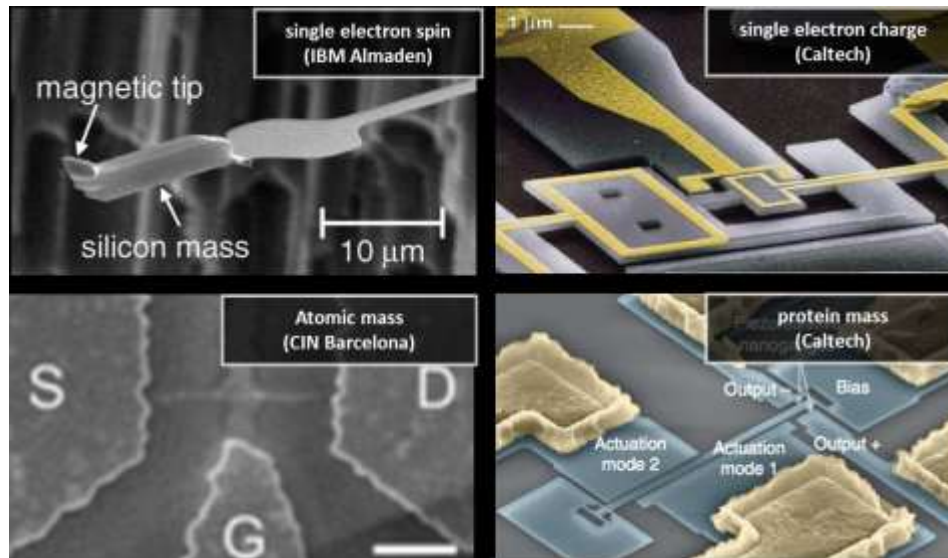
Mechanical quantum sensing

Implementation	Qubit(s)	Measured quantity(ies)	Typical frequency	Implementation	Qubit(s)	Measured quantity(ies)	Typical frequency
Neutral atoms				Superconducting circuits			
Atomic vapor	Atomic spin	Magnetic field, rotation, time/frequency	dc-GHz	SQUID ^c	Supercurrent	Magnetic field	dc-GHz
Cold clouds	Atomic spin	Magnetic field, acceleration, time/frequency	dc-GHz	Flux qubit	Circulating currents	Magnetic field	dc-GHz
				Charge qubit	Charge eigenstates	Electric field	dc-GHz
Trapped ion(s)				Elementary particles			
	Long-lived electronic state	Time/frequency	THz	Muon	Muonic spin	Magnetic field	dc
	Vibrational mode	Rotation		Neutron	Nuclear spin	Magnetic field, phonon density, gravity	dc
		Electric field, force	MHz				
Rydberg atoms				Other sensors			
	Rydberg states	Electric field	dc, GHz	SET ^d	Charge eigenstates	Electric field	dc-MHz
Solid-state spins (ensembles)				Optomechanics	Phonons	Force, acceleration, mass, magnetic field, voltage	kHz–GHz
NMR sensors	Nuclear spins	Magnetic field	dc				
NV ^b center ensembles	Electron spins	Magnetic field, electric field, temperature, pressure, rotation	dc-GHz	Interferometer	Photons, (atoms, molecules)	Displacement, refractive index	...
Solid-state spins (single spins)							
P donor in Si	Electron spin	Magnetic field	dc-GHz				
Semiconductor quantum dots	Electron spin	Magnetic field, electric field	dc-GHz				
Single NV ^b center	Electron spin	Magnetic field, electric field, temperature, pressure, rotation	dc-GHz				

* C. L. Degen *et.al*, “Quantum sensing”, *Rev. Mod. Phys.* **89**, 035002 (2017).

Mechanical oscillators as force sensors

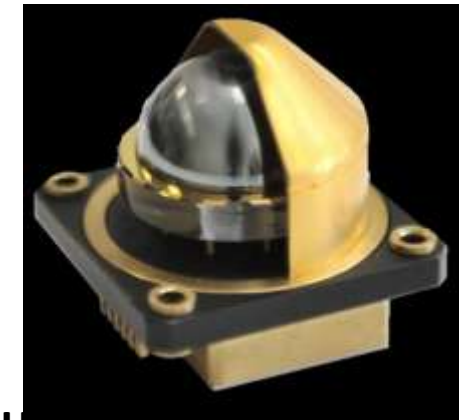
Scientific



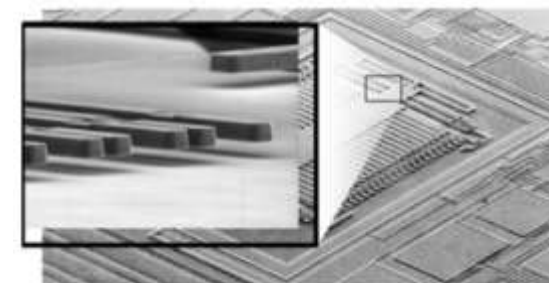
MEMS/NEMS sensors

- Force from single quanta
- Mass of single atom/molecule

Industrial

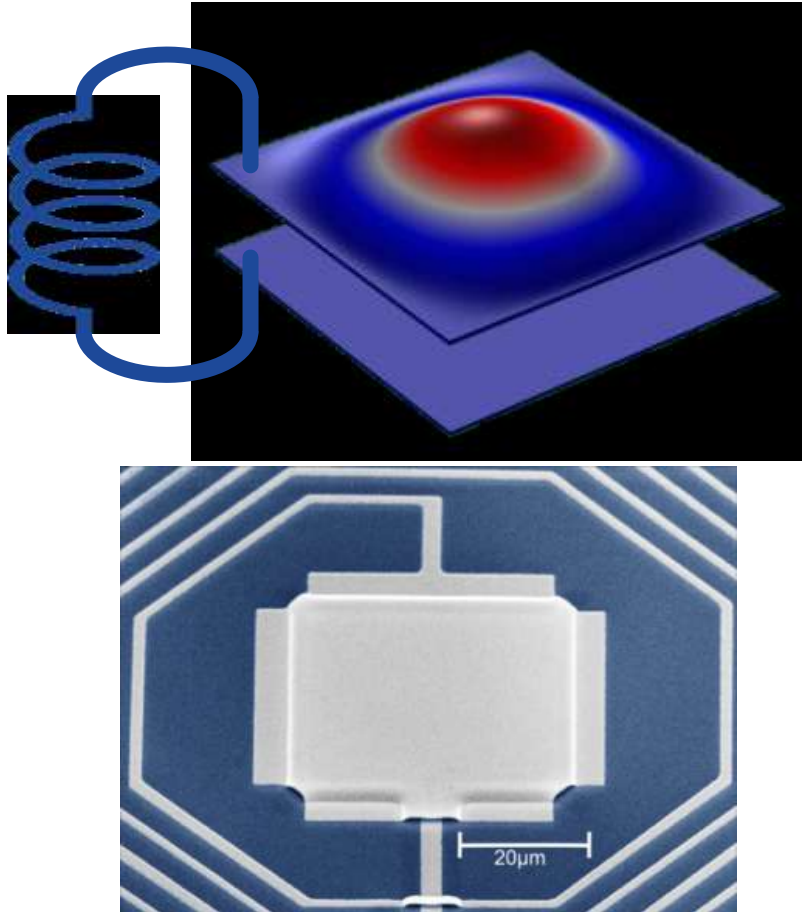


Hemispherical resonator gyroscope
- Coriolis force

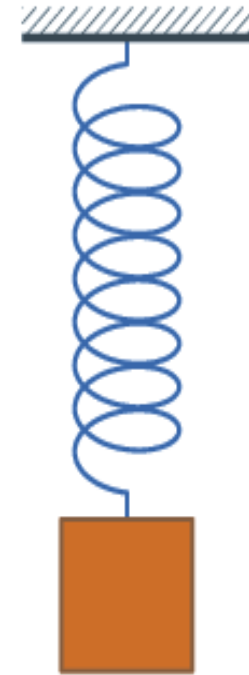
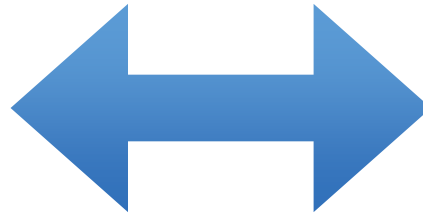


MEMS accelerometer
- inertial force from acceleration

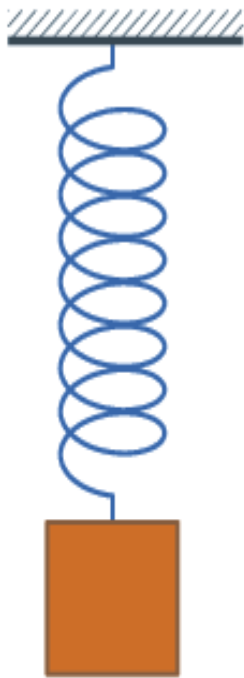
역학적 진동자



* Suh *et al.*, *Science* **344**, 1262 (2014)

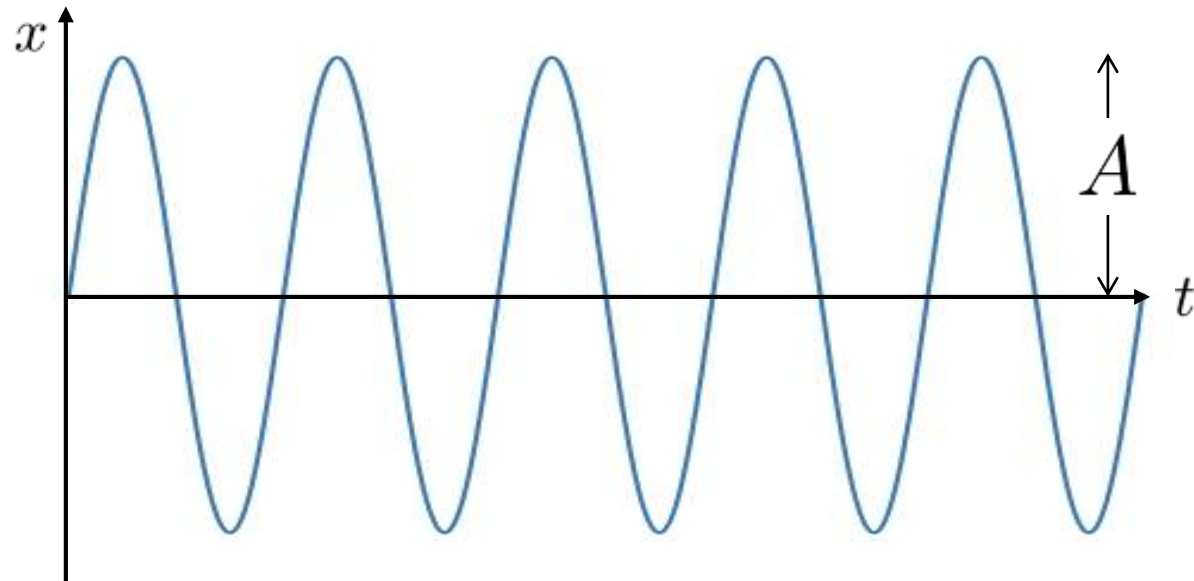


조화 진동자

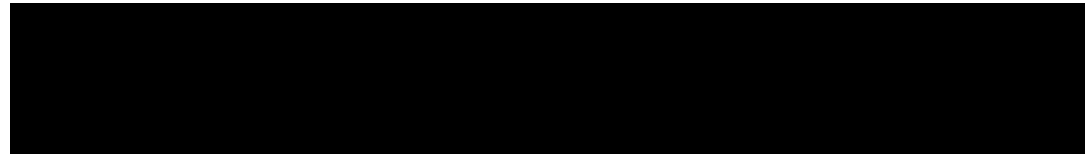
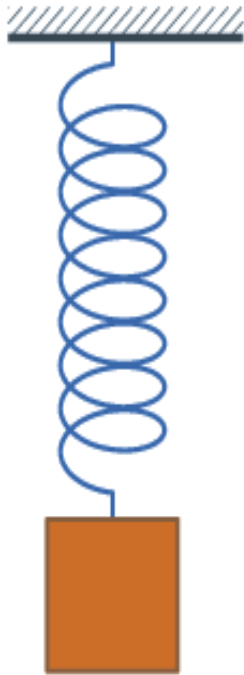


for external force $F \cos \omega t$,

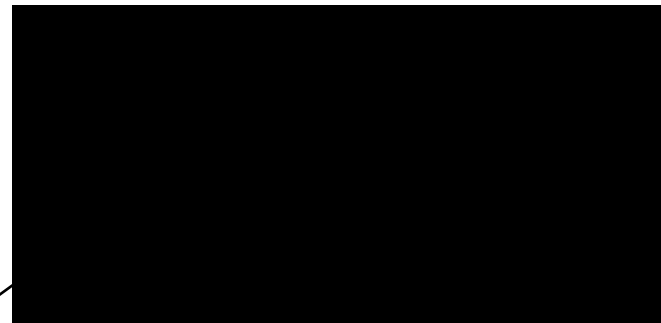
$$m \frac{d^2 x}{dt^2} + \frac{m \omega_0}{Q} \frac{dx}{dt} + m \omega_0^2 x = F \cos \omega t$$



Amplitude measurement



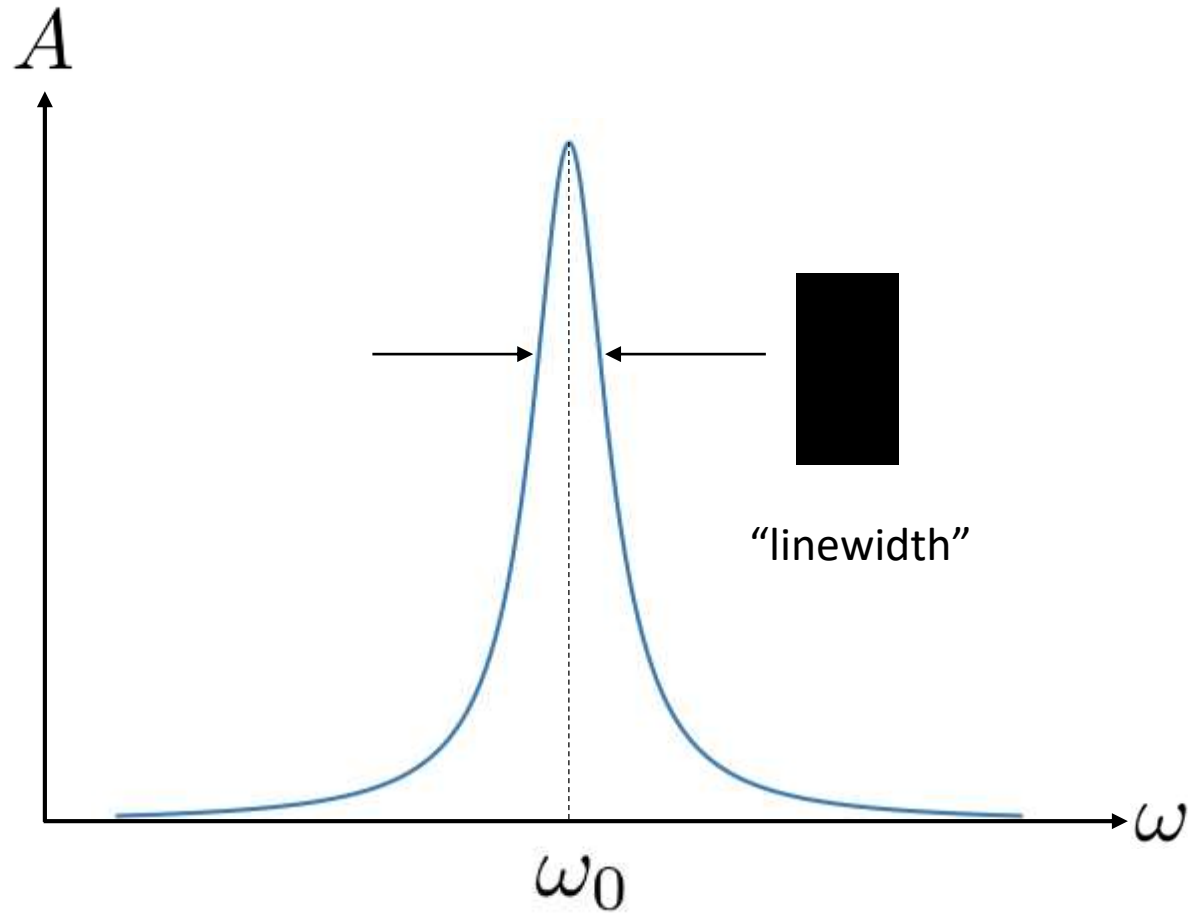
측정량



Q-값 (품질값); $10 \sim 10^9$

스프링계수; $10^{-4} \sim 10^2$ N/m

Near-resonant Force



Frequency measurement

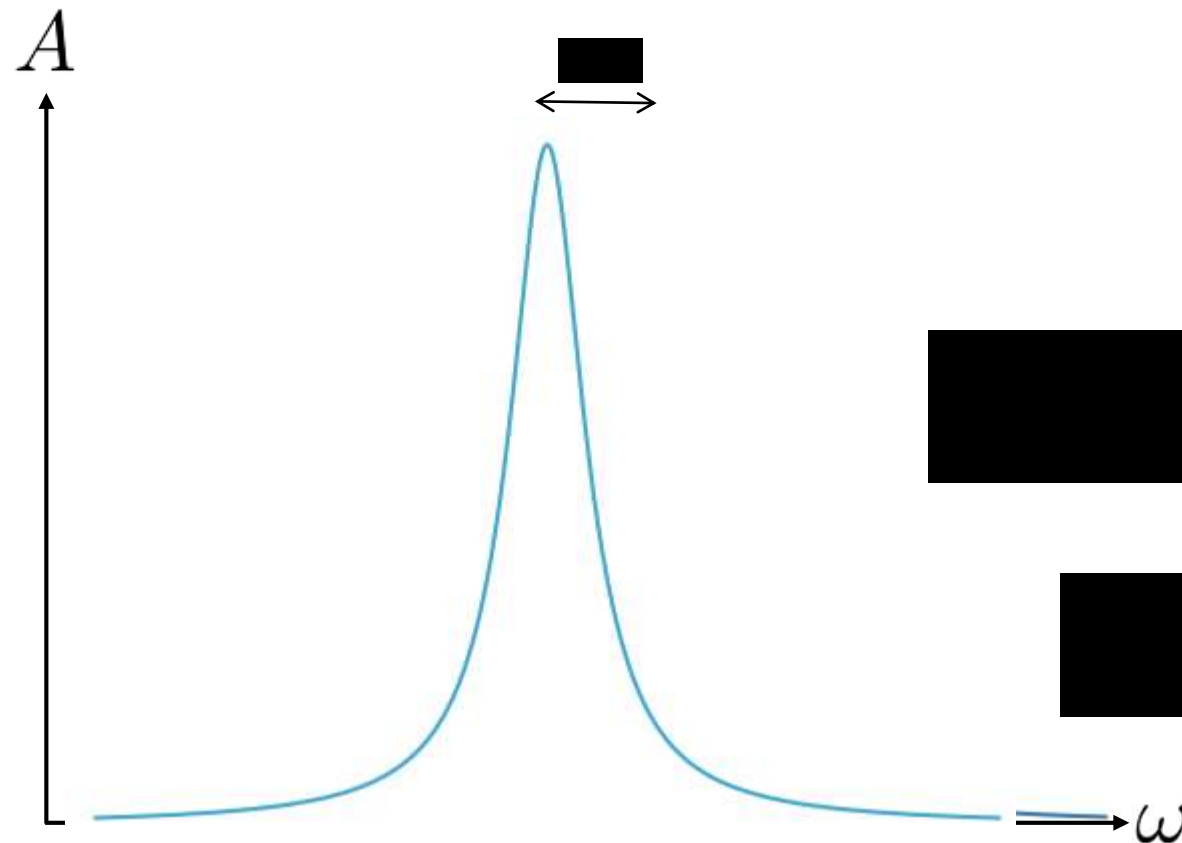
$$\omega_0 = \sqrt{\frac{k}{m}}$$

측정량 ω_0

스프링계수; $10^{-4} \sim 10^2 \text{ N/m}$

질량; $\text{pg} \sim \text{ng}$

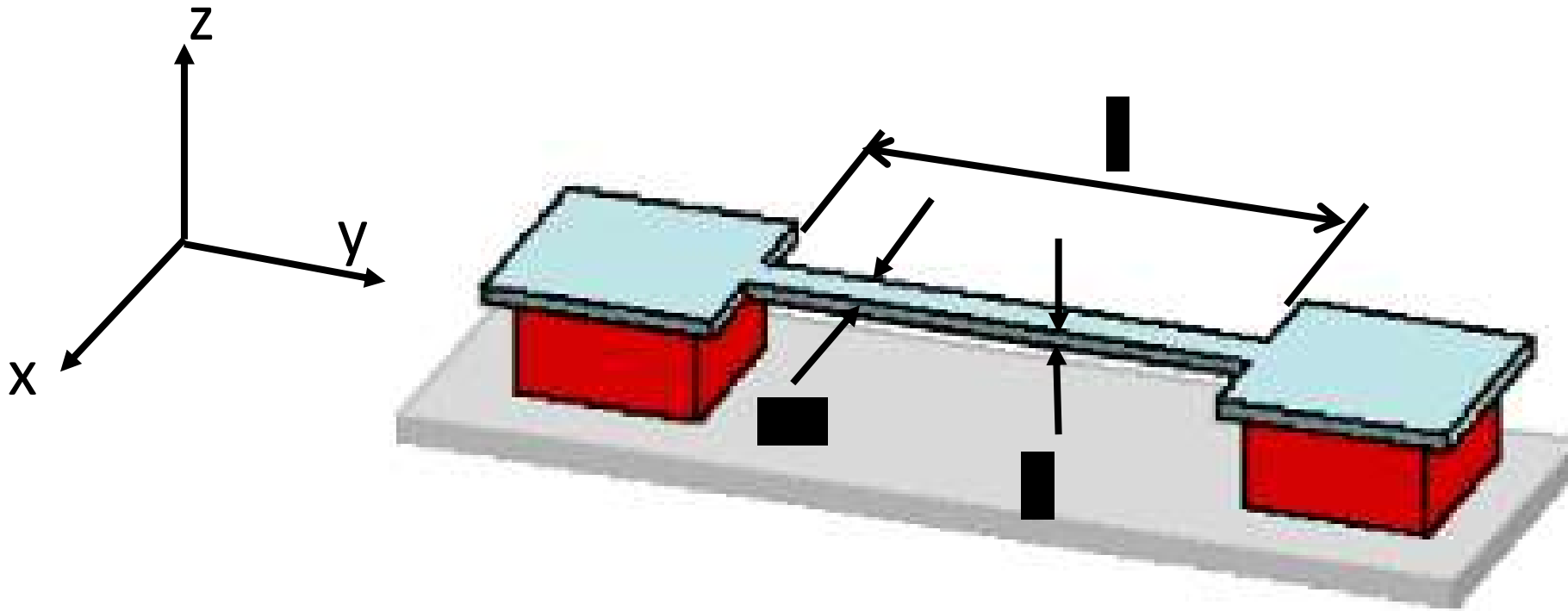
Frequency measurement



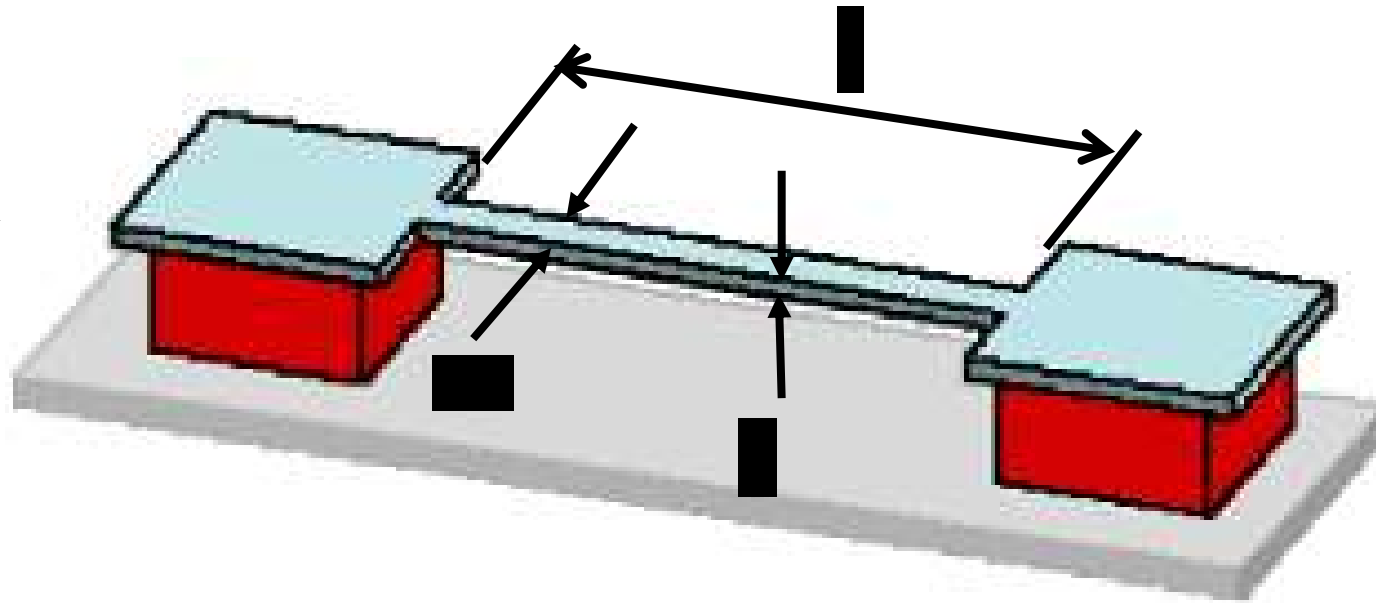
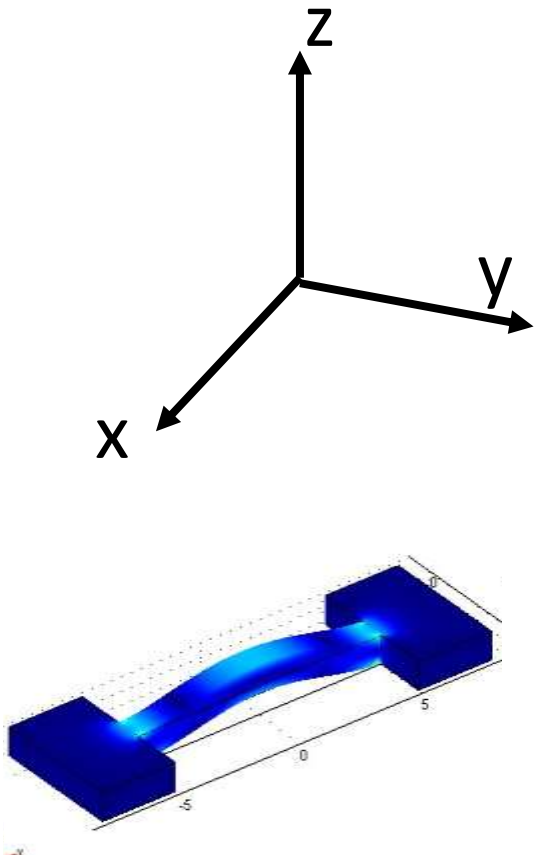
스프링계수 or
force gradient
(k)의 변화

질량의 변화

Example of doubly clamped beam

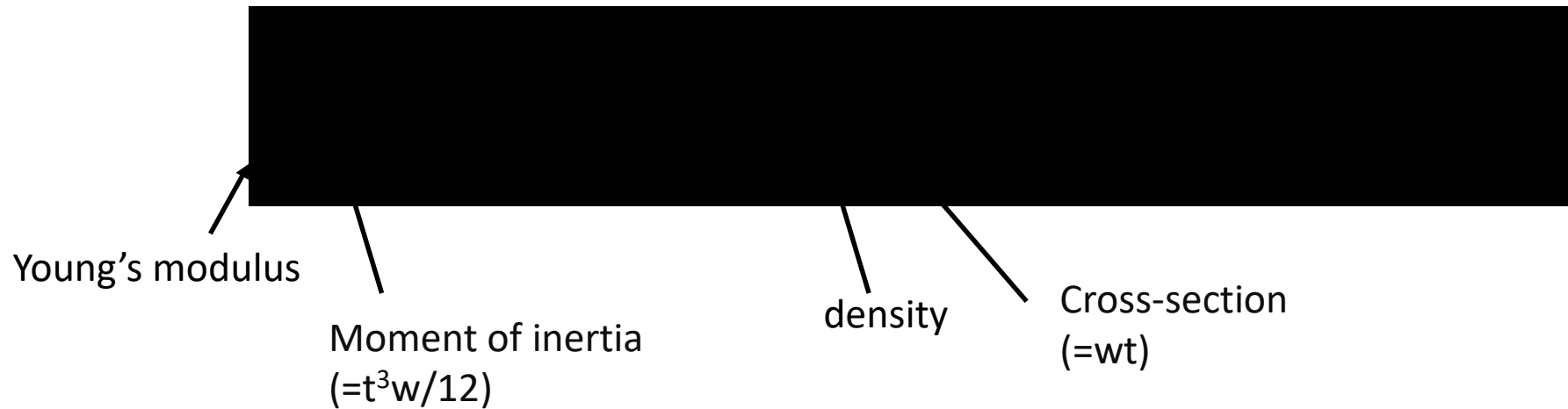


Eigenmode = harmonic oscillator



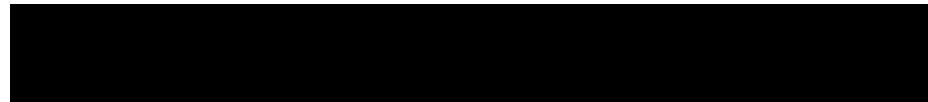
Equation of motion

- Euler-Bernulli equation



The diagram shows a large black rectangular box representing the Euler-Bernulli equation. Four lines point from labels below to different parts of the box: 'Young's modulus' points to the left side, 'Moment of inertia (=t³w/12)' points to the bottom-left corner, 'density' points to the bottom-right corner, and 'Cross-section (=wt)' points to the right side.

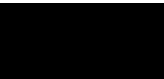
- seperation of variables; normal modes



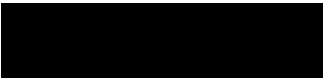
A solid black rectangular box representing the equation for separation of variables and normal modes.

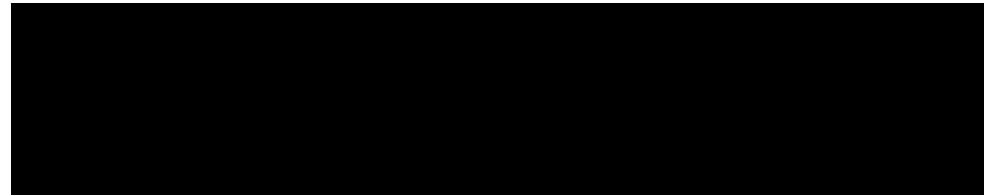
* Foundations of nanomechanics, A. N. Cleland

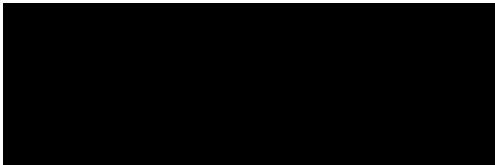
Equation of motion

- damping can be included by adding a dissipation term 



- In frequency domain (quality factor ),

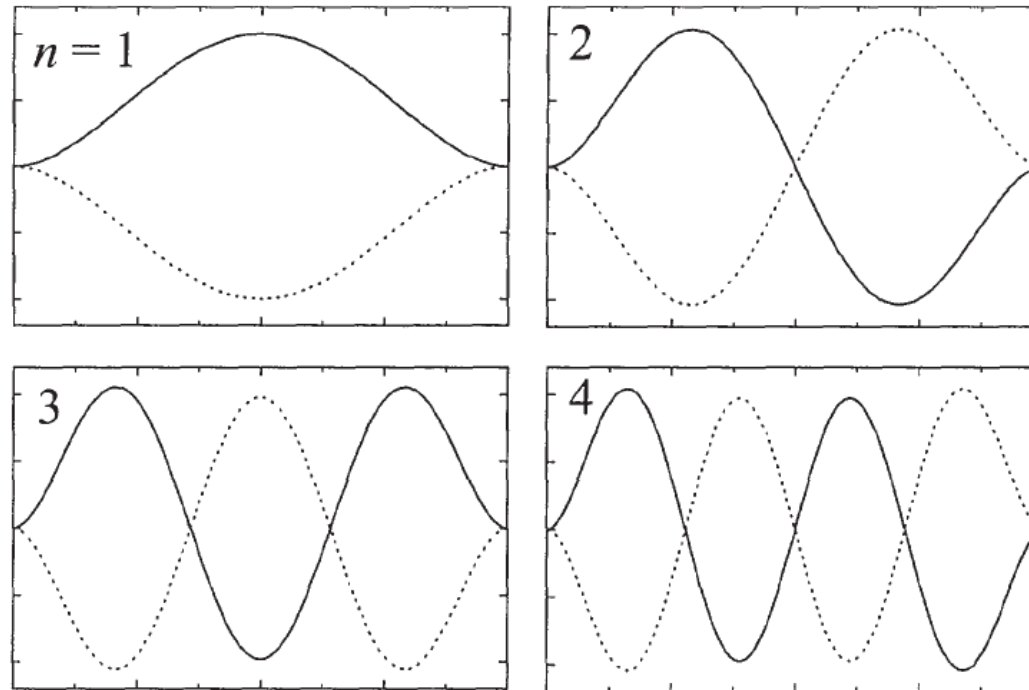
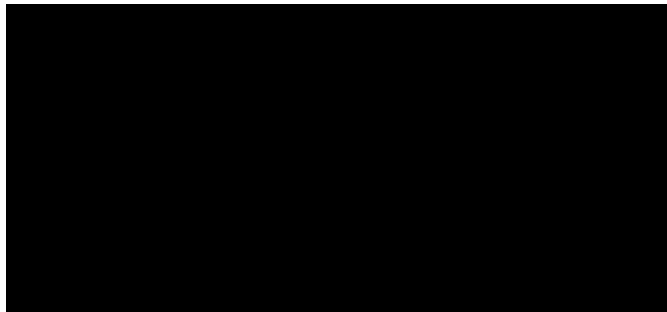


- At resonance, 

* Foundations of nanomechanics, A. N. Cleland

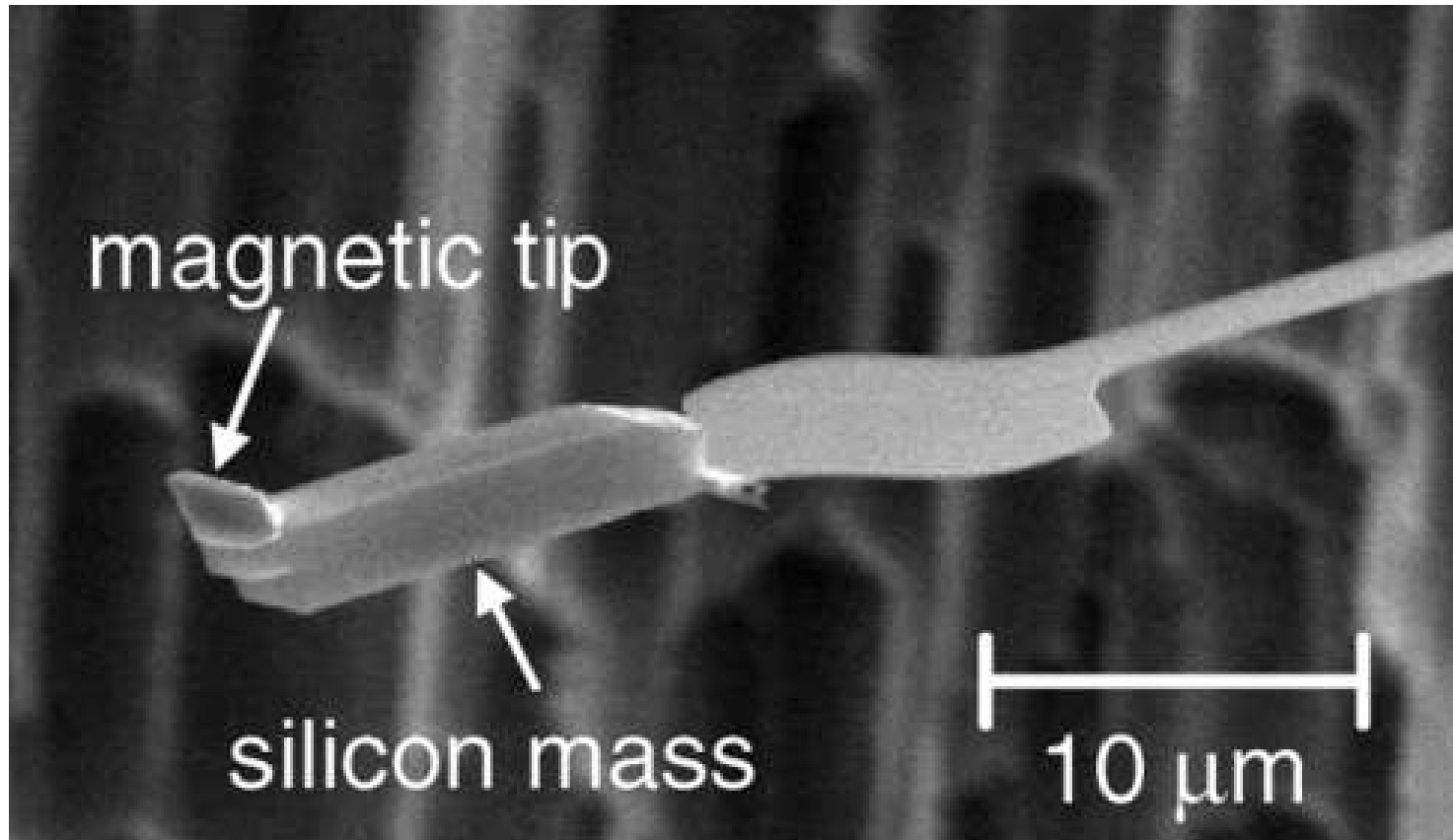
Example of doubly-clamped beam

- boundary condition: fixed ends, zero-slopes
- First four mode shapes



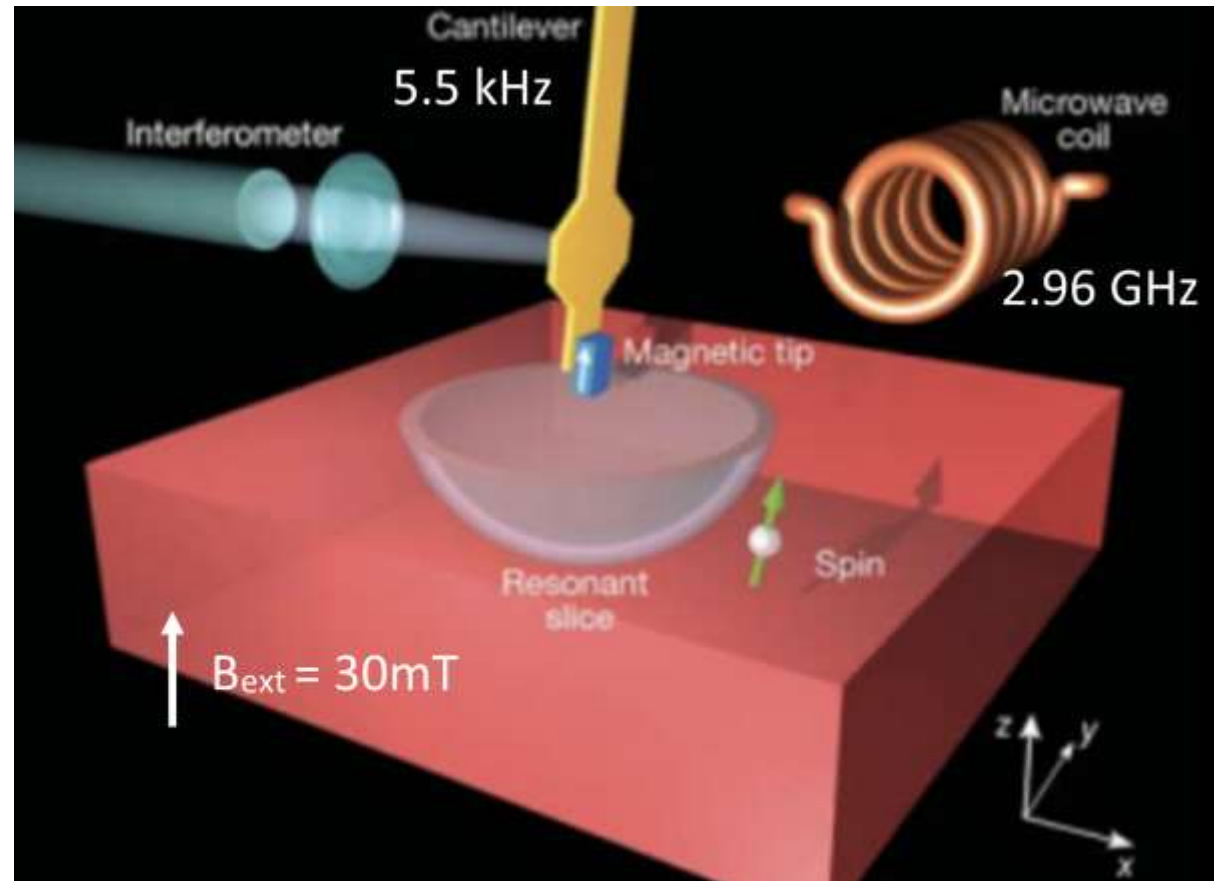
* Foundations of nanomechanics, A. N. Cleland

예제: magnetic resonance force microscopy



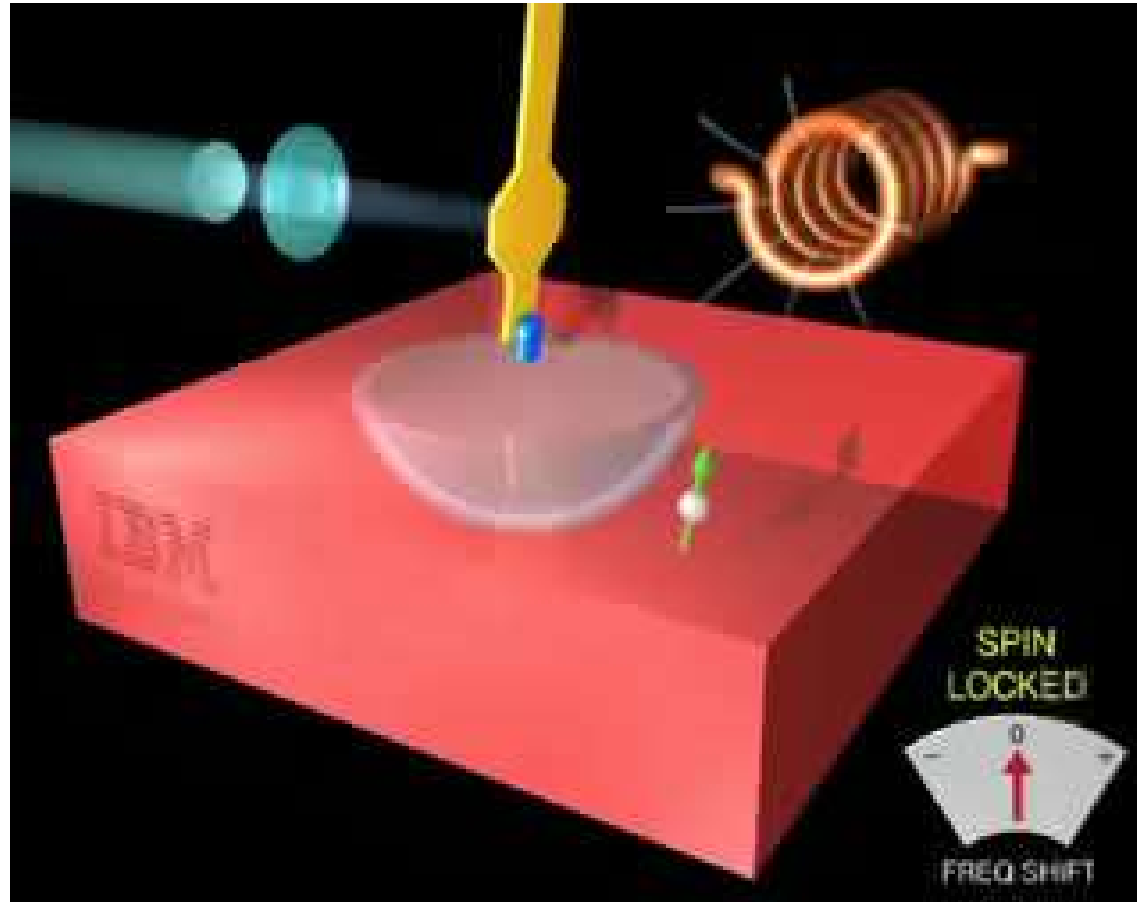
* D. Rugar *et.al.*, Single spin detection by magnetic resonance force microscopy, *Nature* **430**, 329 (2004).

Mechanical detection of single electron spin flip



* D. Rugar *et.al.*, Single spin detection by magnetic resonance force microscopy, *Nature* **430**, 329 (2004).

Mechanical detection of single electron spin flip

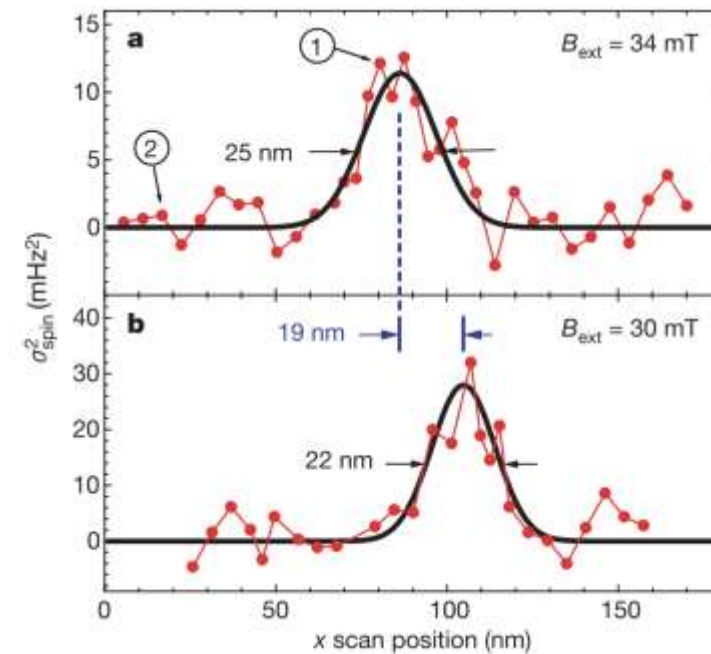
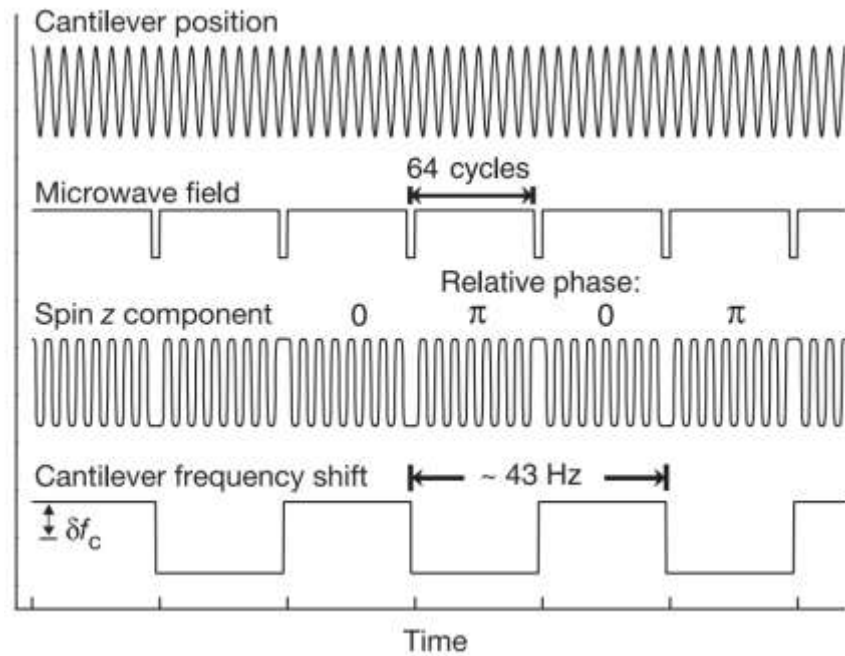


* D. Rugar *et.al.*, Single spin detection by magnetic resonance force microscopy, *Nature* **430**, 329 (2004).

Frequency measurement

$$\delta f_c = \pm \frac{2f_c G \mu_B}{\pi k x_{\text{peak}}}$$

$$G \equiv \partial B_0 / \partial x$$



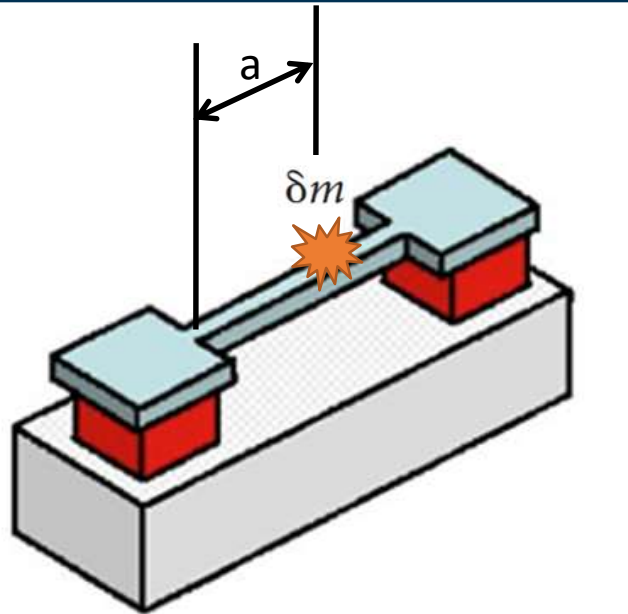
* D. Rugar *et.al.*, Single spin detection by magnetic resonance force microscopy, *Nature* **430**, 329 (2004).

예제: protein mass spectrometry

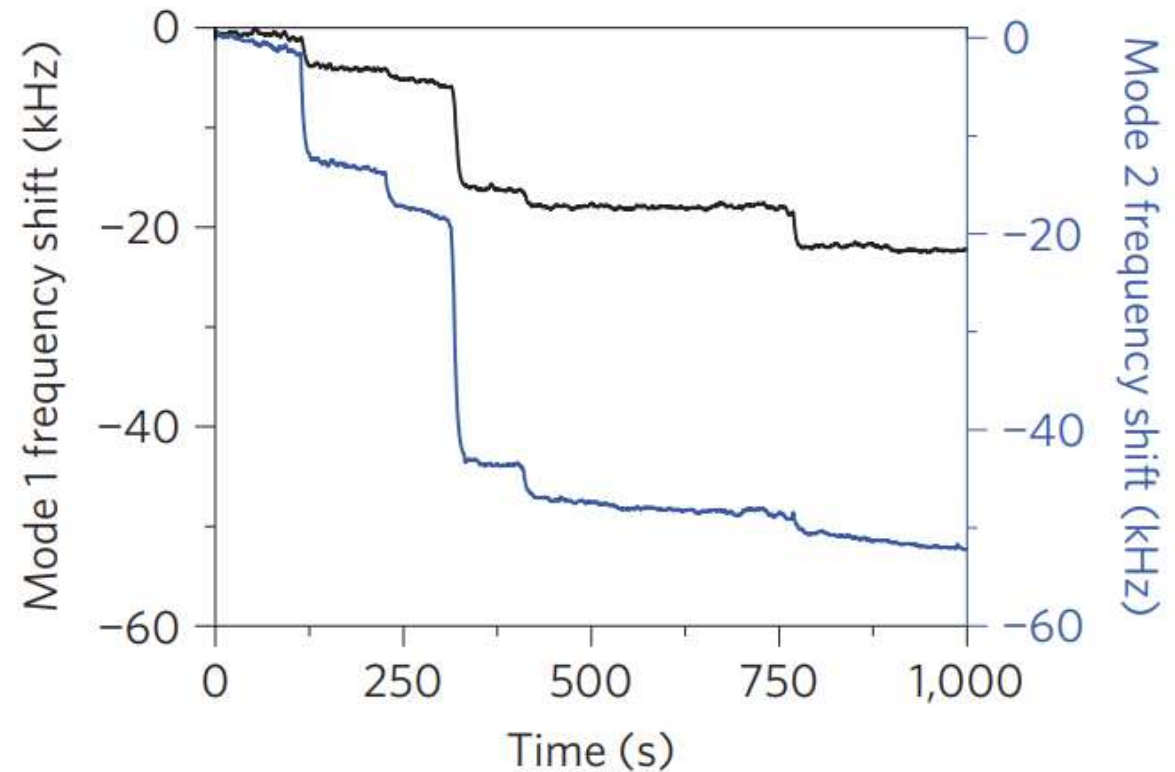


* M. S. Hanay *et.al.*, Single-protein nanomechanical mass spectrometry in real time, *Nat. Nano.* **7**, 602 (2012).

Mechanical mass sensing

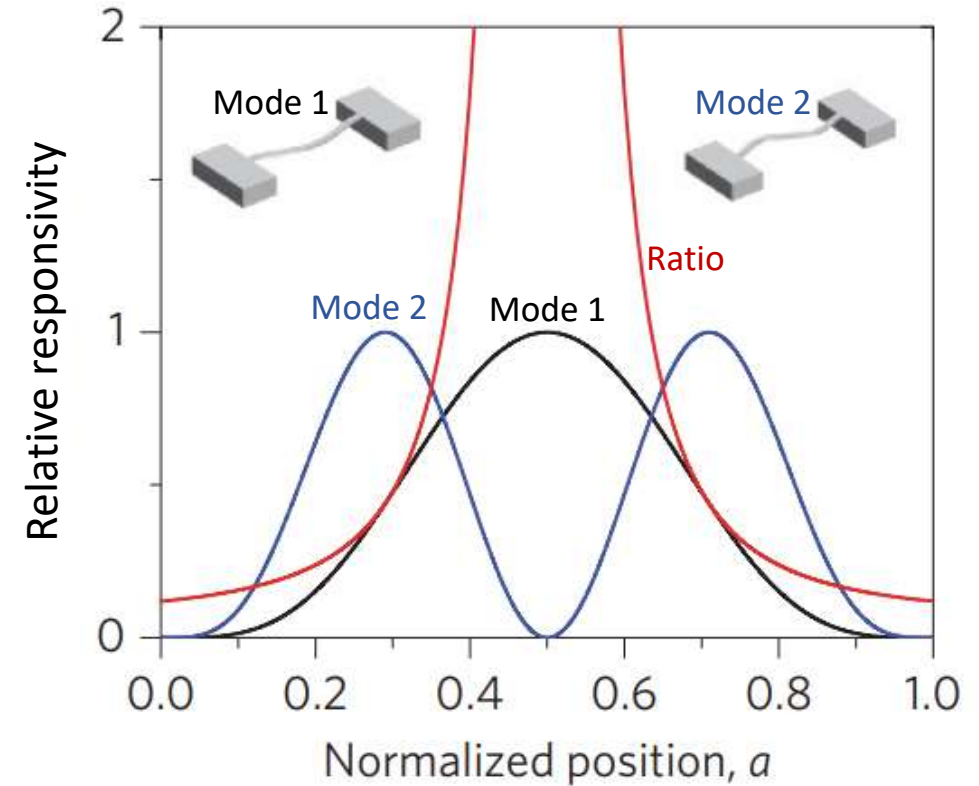
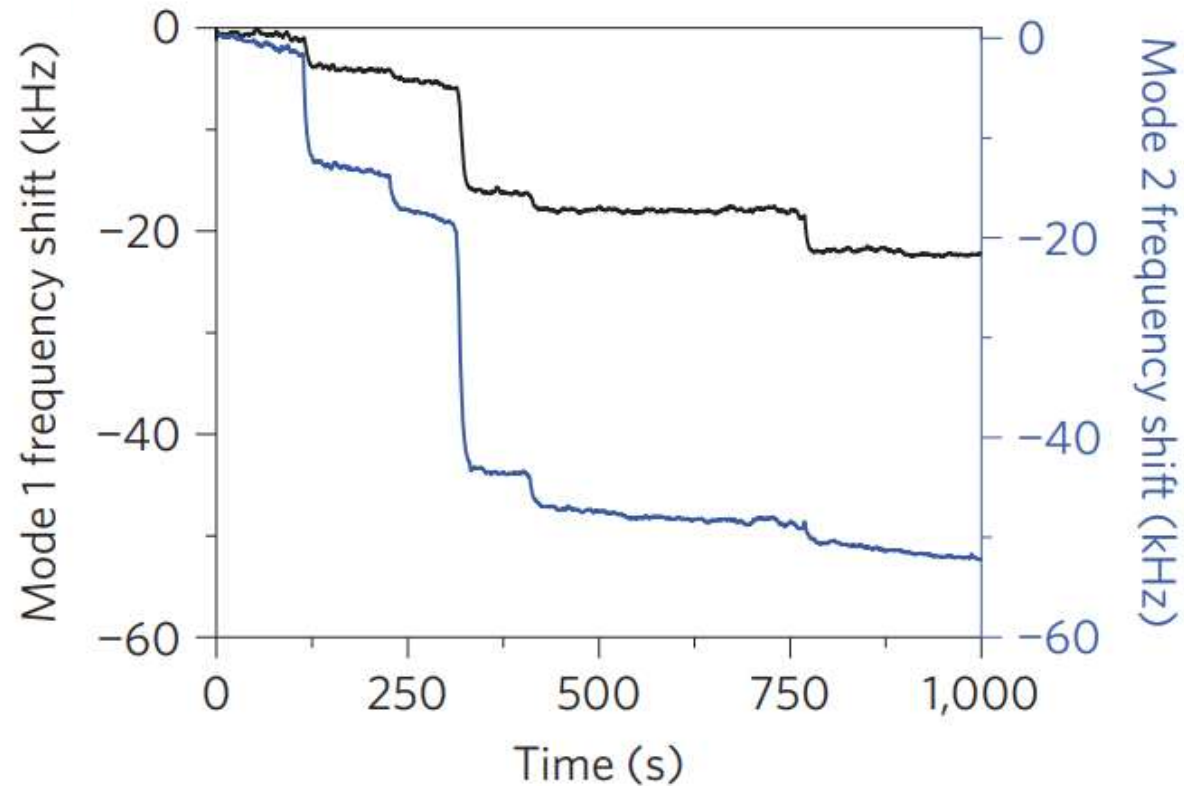


$$\frac{\Delta f_n}{f_n} = -\frac{\delta m \psi_n(a)^2}{M \alpha_n^2}$$



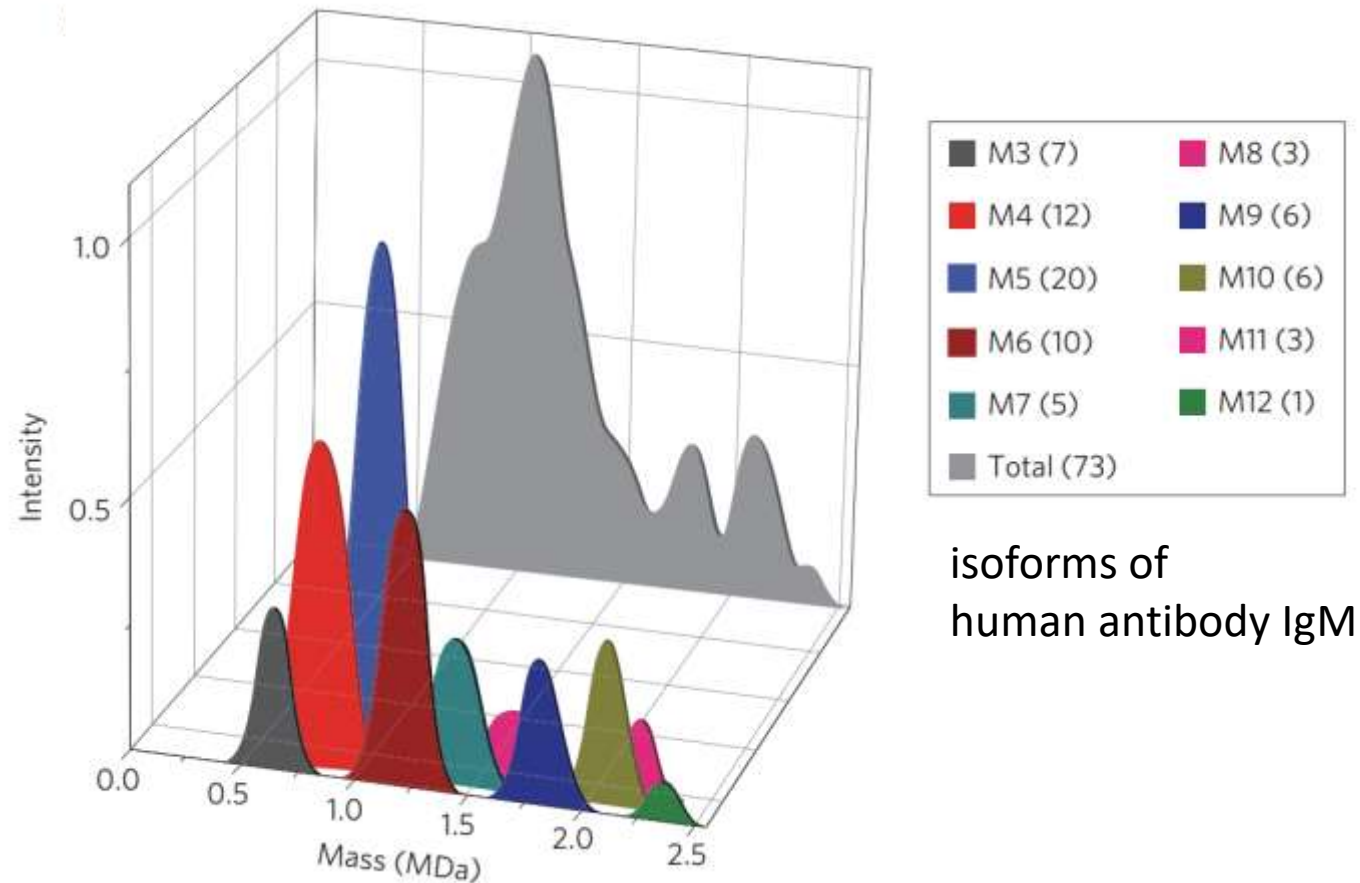
* M. S. Hanay *et.al.*, Single-protein nanomechanical mass spectrometry in real time, *Nat. Nano.* **7**, 602 (2012).

Multi-mode frequency measurement



* M. S. Hanay *et.al.*, Single-protein nanomechanical mass spectrometry in real time, *Nat. Nano.* **7**, 602 (2012).

Protein mass spectrometry



* M. S. Hanay *et.al.*, Single-protein nanomechanical mass spectrometry in real time, *Nat. Nano.* **7**, 602 (2012).

Quantum mechanics defines minimum uncertainty in position measurement



“standard quantum limit”

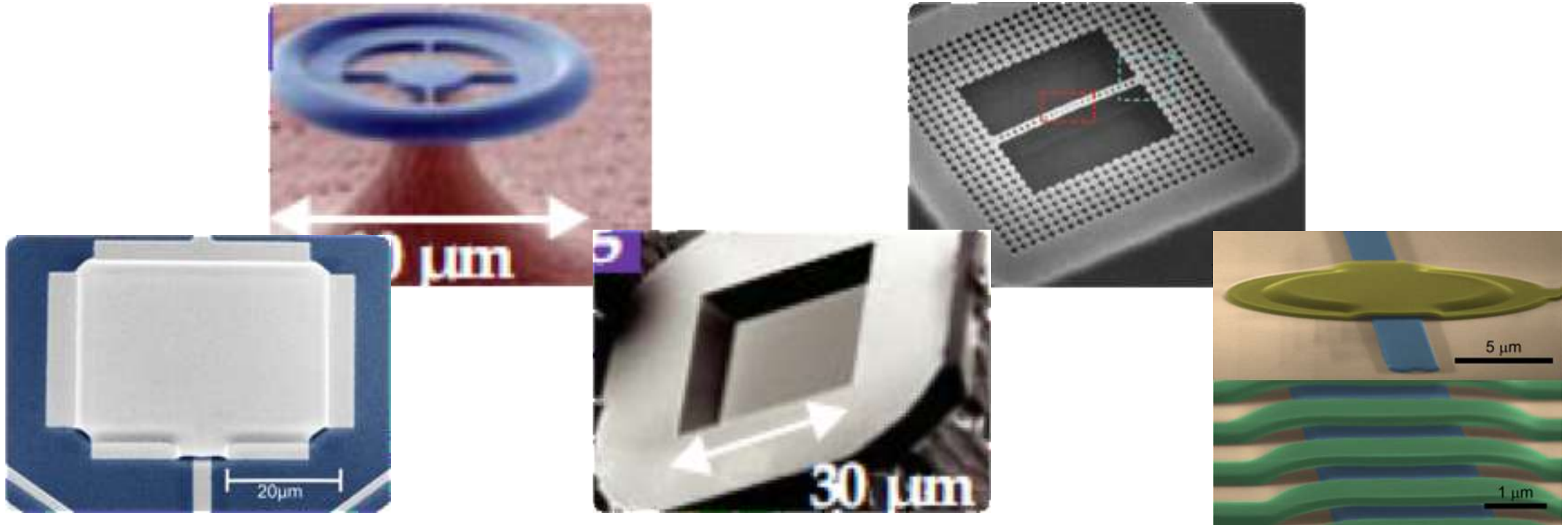
* C. M. Caves *et.al.*, *Rev. Mod. Phys.* **52**, 341 (1980).

Quantum mechanics defines
minimum uncertainty in position measurement

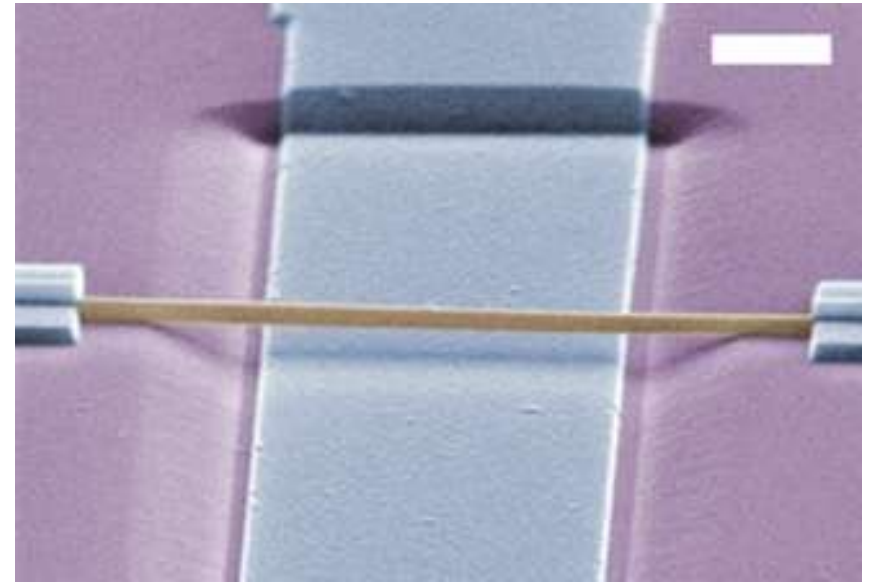
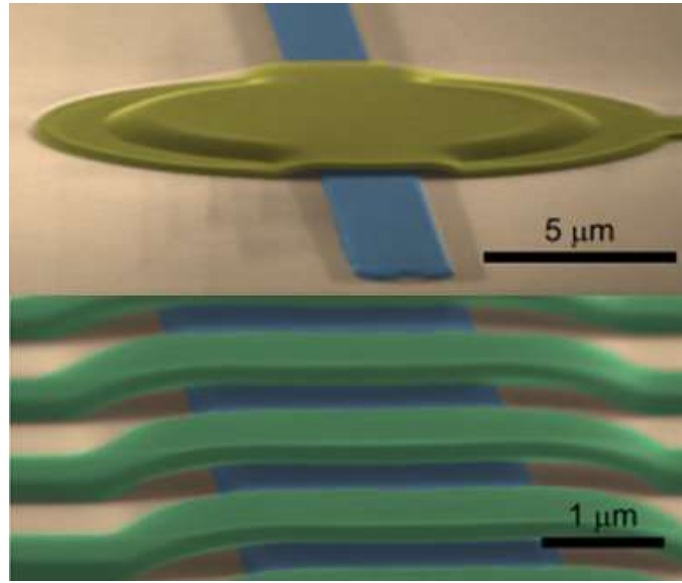
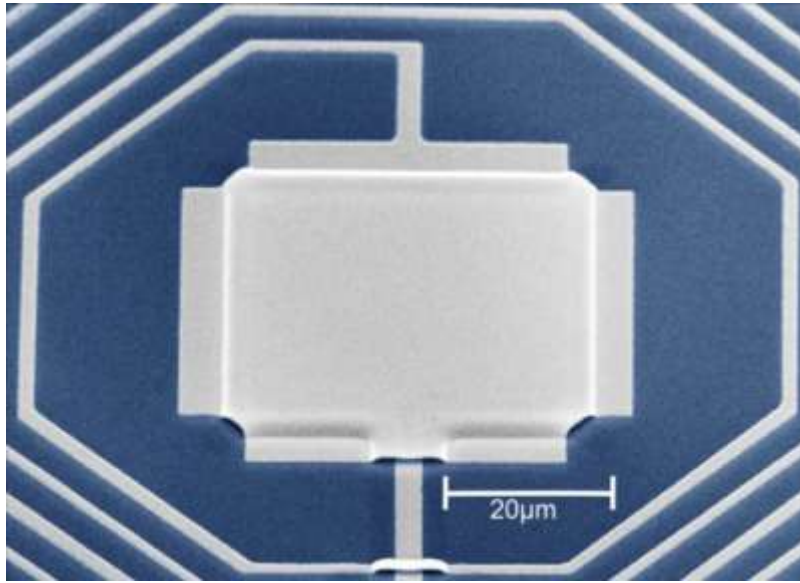
Quantum measurement of mechanical oscillators II Ultimate-precision force sensing

* C. M. Caves *et.al.*, *Rev. Mod. Phys.* **52**, 341 (1980).

Cavity optomechanical system (광역학계)



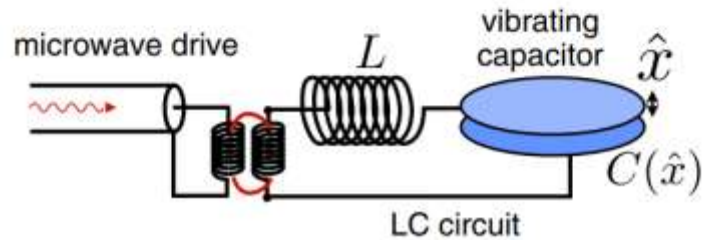
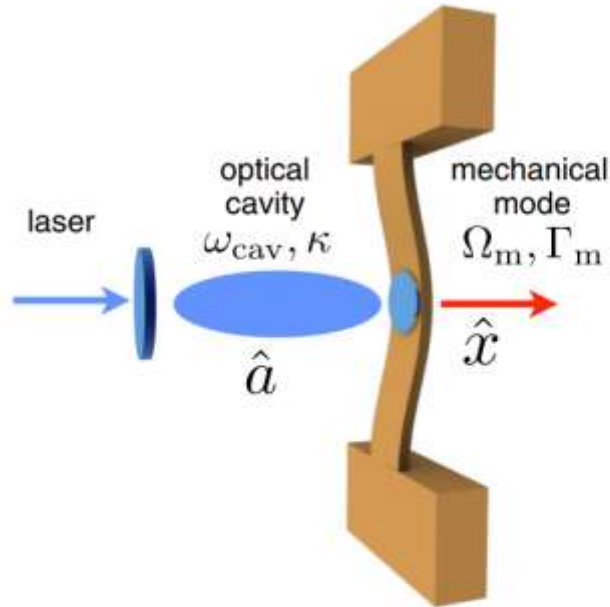
Strong coupling between phonons and photons
⇒ *hybrid quantum devices for quantum information science*



중시계 진동자의 광역학적 상호작용

Optomechanical interactions in mesoscopic oscillators

Cavity optomechanical system (광역학계)



[Cavity Hamiltonian]

$$= \hbar \omega_{cav} \hat{a}^\dagger \hat{a}$$



[Optomechanical interaction Hamiltonian]

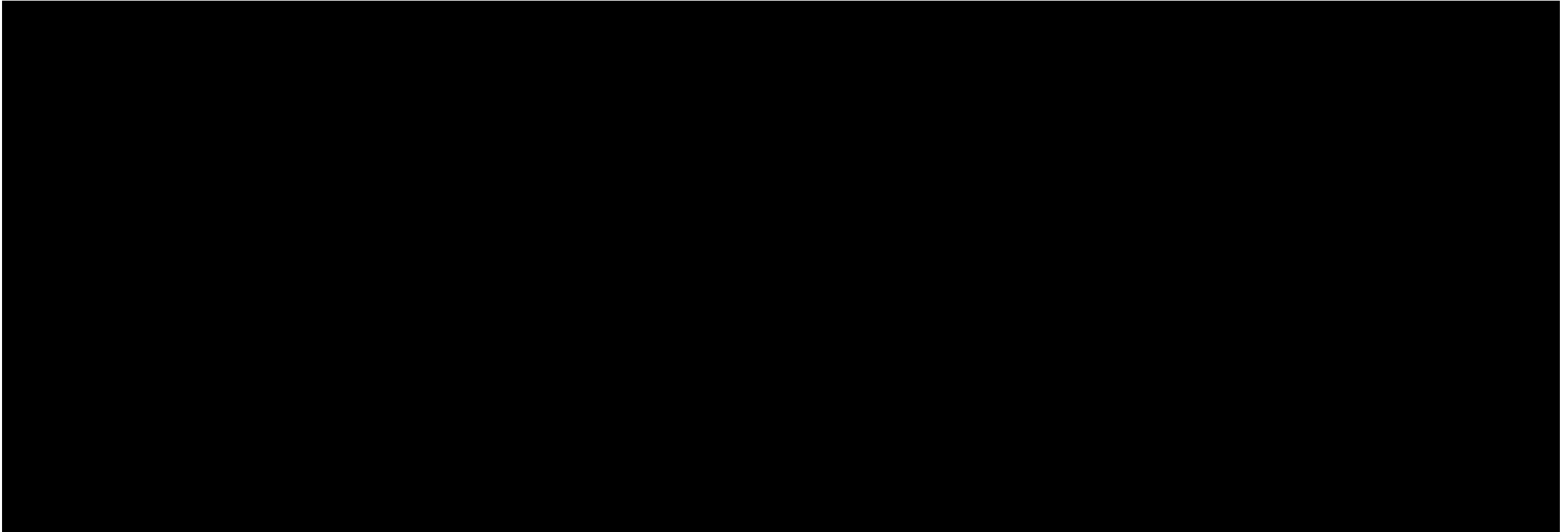
$$= \hbar \frac{\partial \omega_{cav}}{\partial x} \hat{x} \hat{a}^\dagger \hat{a} \quad \text{"radiation pressure"}$$

$$\begin{aligned} \Rightarrow L_{cav} = n\lambda = n \frac{2\pi c}{\omega_{cav}} &\rightarrow \delta x = \delta L_{cav} \approx -n \frac{2\pi c}{\omega_{cav}^2} \delta \omega_{cav} \\ \frac{\partial \omega_{cav}}{\partial x} &\approx -\frac{\omega_{cav}^2}{2\pi c n} \end{aligned}$$

$$\begin{aligned} \Rightarrow \omega_{cav} = \frac{1}{\sqrt{LC}}; C = \frac{\epsilon_0 A_{cap}}{d_{cap}} &\rightarrow \delta x = \delta d_{cap} \approx \frac{2d_{cap}}{\omega_{cav}} \delta \omega_{cav} \\ \frac{\partial \omega_{cav}}{\partial x} &\approx \frac{\omega_{cav}}{2d_{cap}} \end{aligned}$$

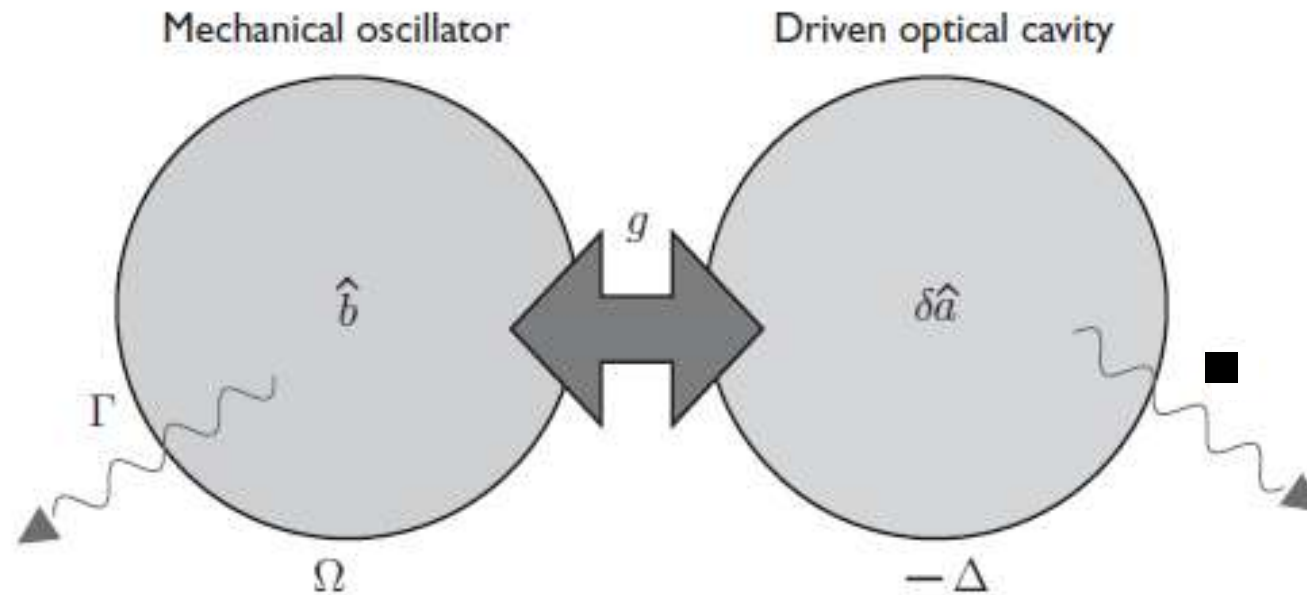
* Aspelmeyer et al., Rev. Mod. Phys. **86**, 29 (2014).

Optomechanical interaction



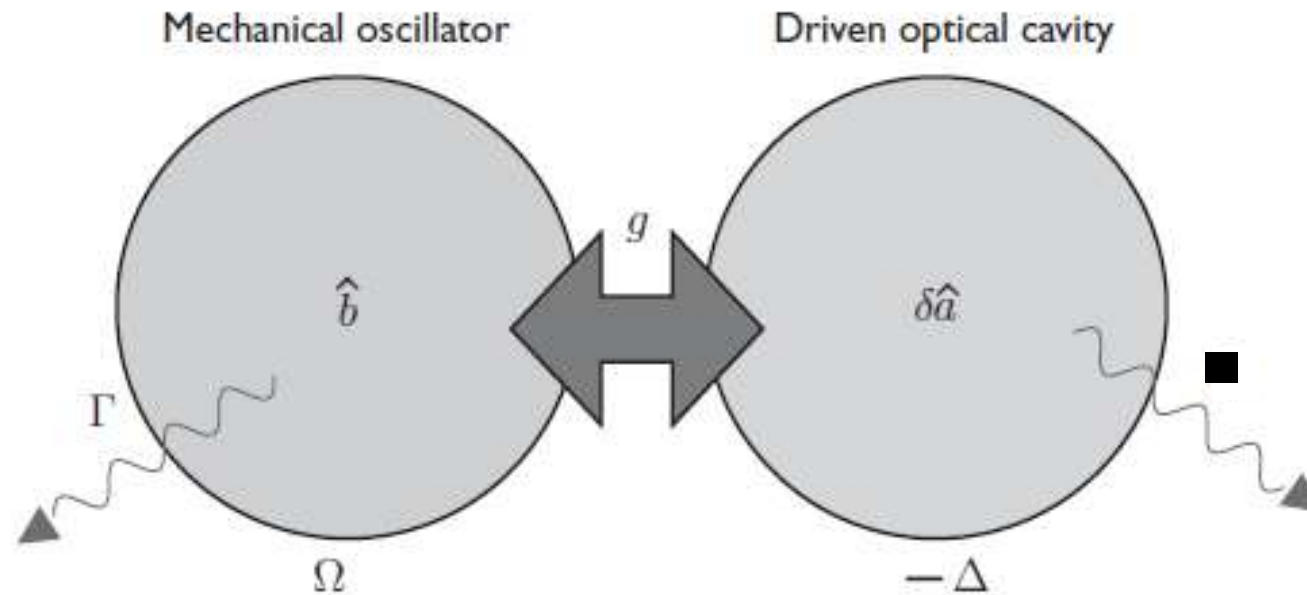
* Devoret *et.al.*, *lecture notes of les houches summer school* (2011).

Optomechanical interaction



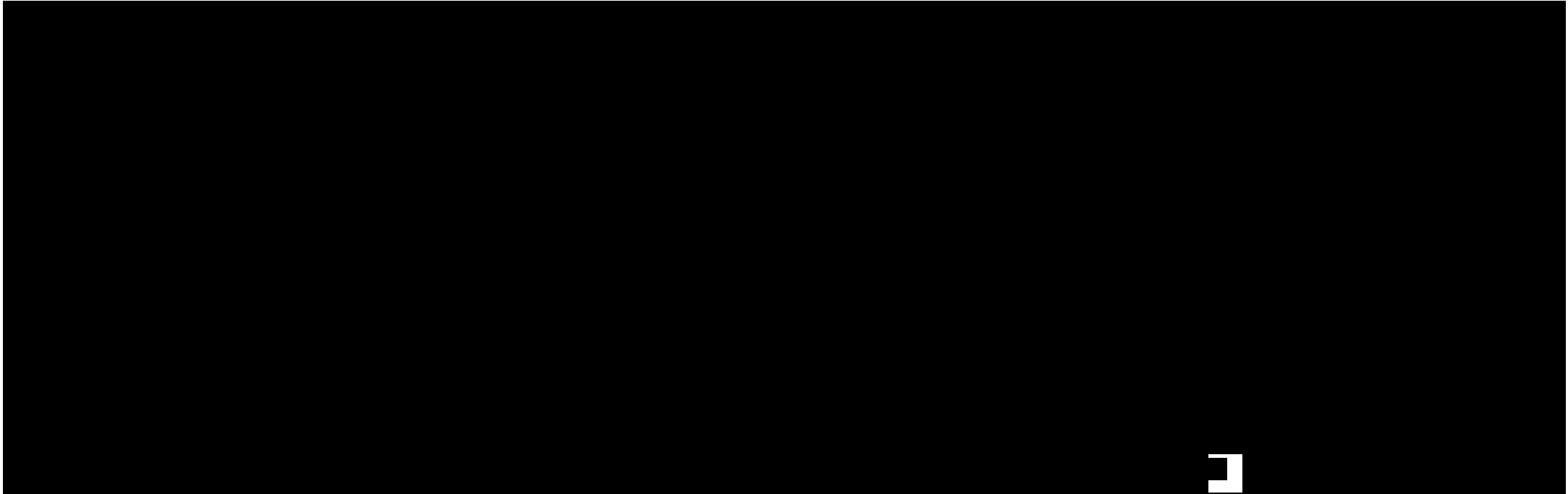
* Devoret *et.al.*, *lecture notes of les houches summer school* (2011).

Optomechanical interaction



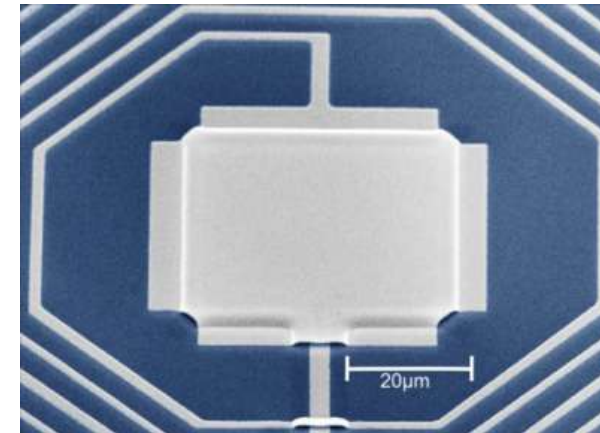
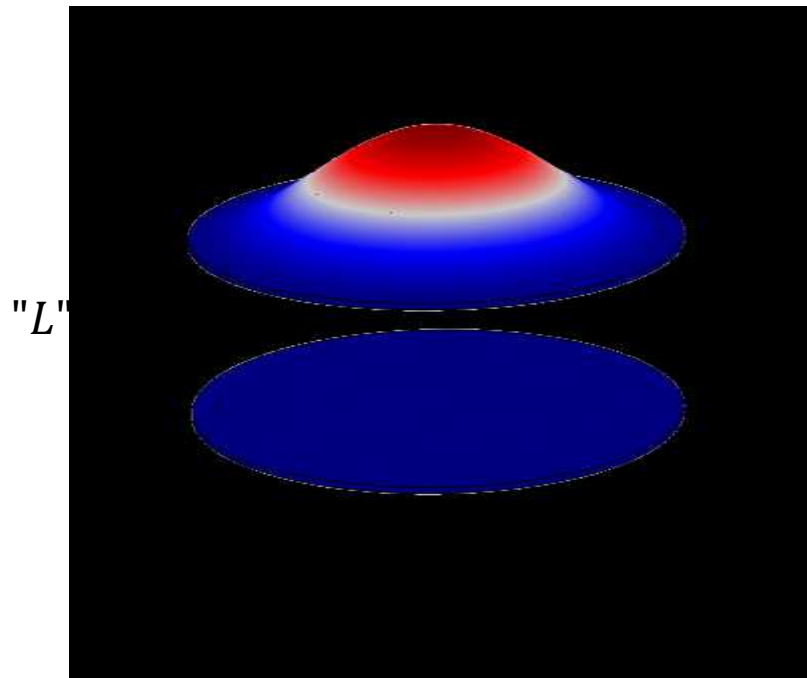
* Devoret *et.al.*, *lecture notes of les houches summer school* (2011).

Optomechanical interaction



* Devoret *et.al.*, *lecture notes of les houches summer school* (2011).

Microwave cavity optomechanical system



* JS et.al., *Science* **344**, 1262 (2014).

$$\hbar\omega_m \hat{b}^\dagger \hat{b} + \hbar g \hat{a}^\dagger \hat{a} (\hat{b}^\dagger + \hat{b})$$

Photon

$$\omega_c = 5.4 \text{ GHz}$$

$$\kappa = 0.9 \text{ MHz}$$

Phonon

$$\omega_m = 4 \text{ MHz}$$

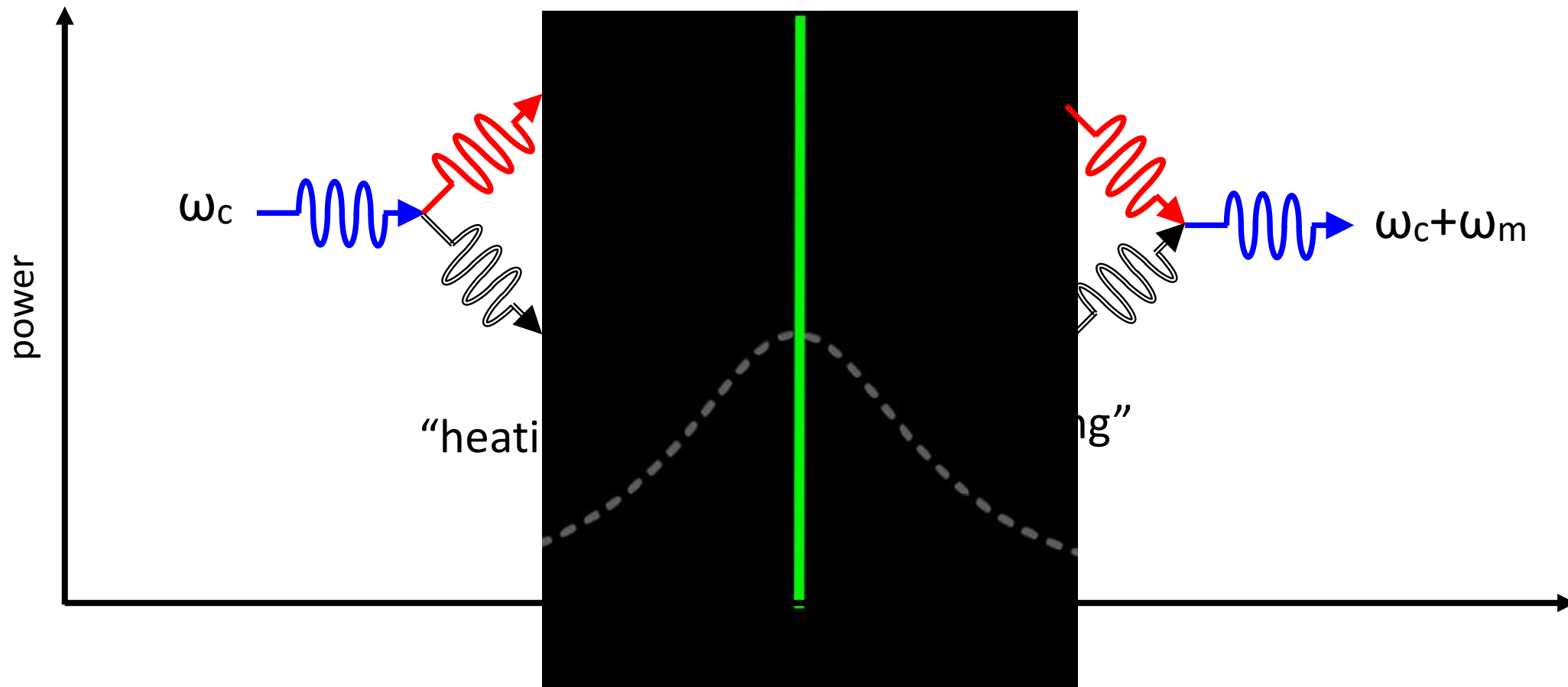
$$\Gamma_m = 10 \text{ Hz}$$

$$x_{zp} = 2 \text{ fm}$$

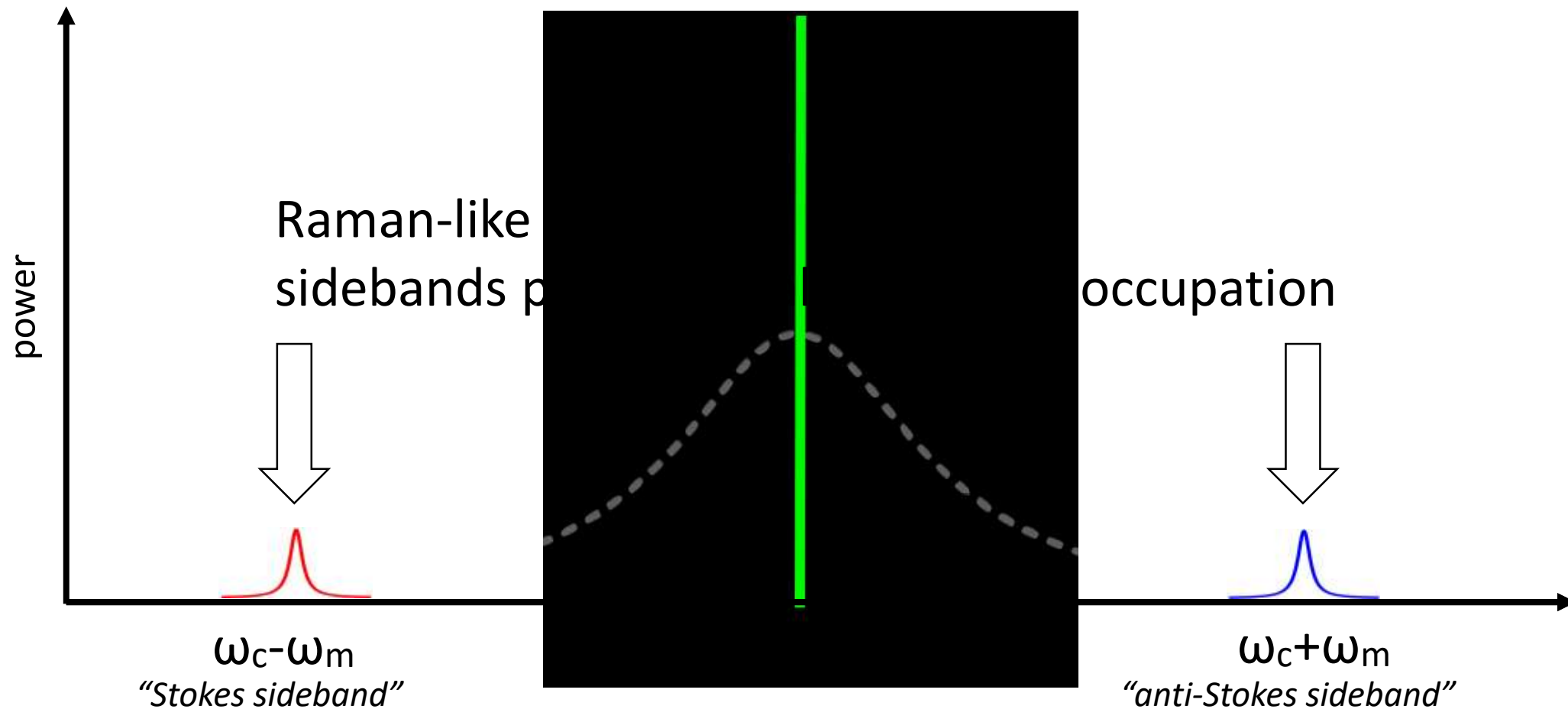
Interaction

$$g = 14 \text{ Hz}$$

Photon-phonon coupling

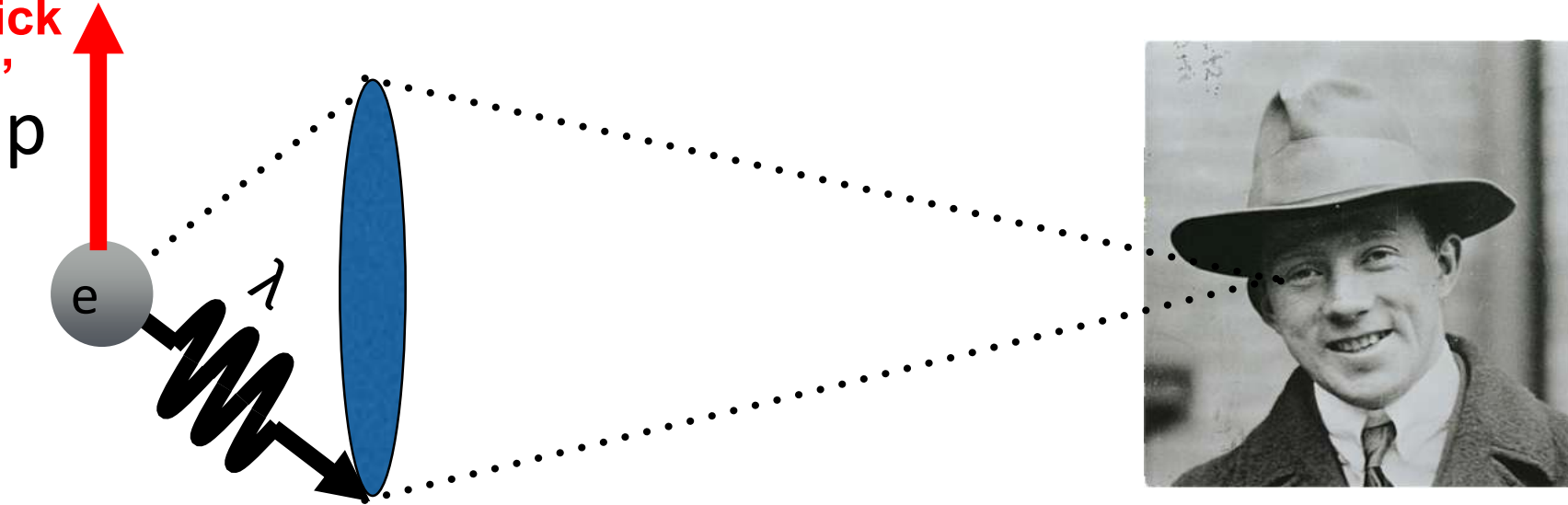


Quantum-limited detection of motion

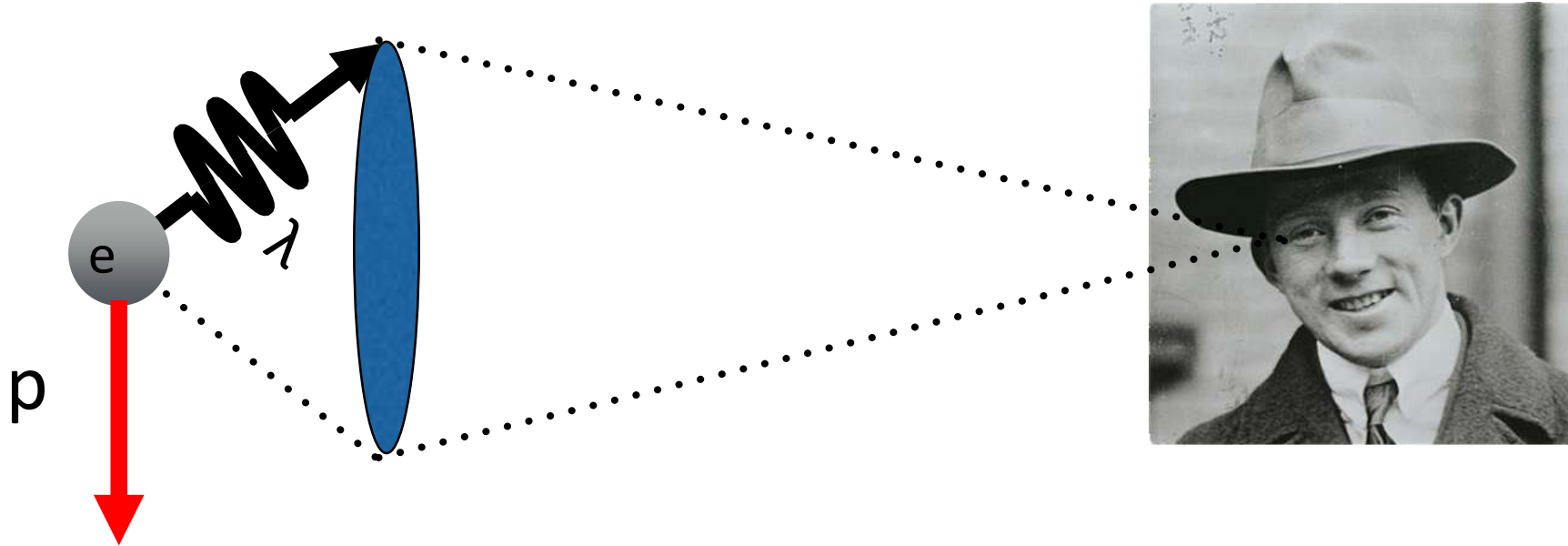


Measurement back-action in Heisenberg's microscope

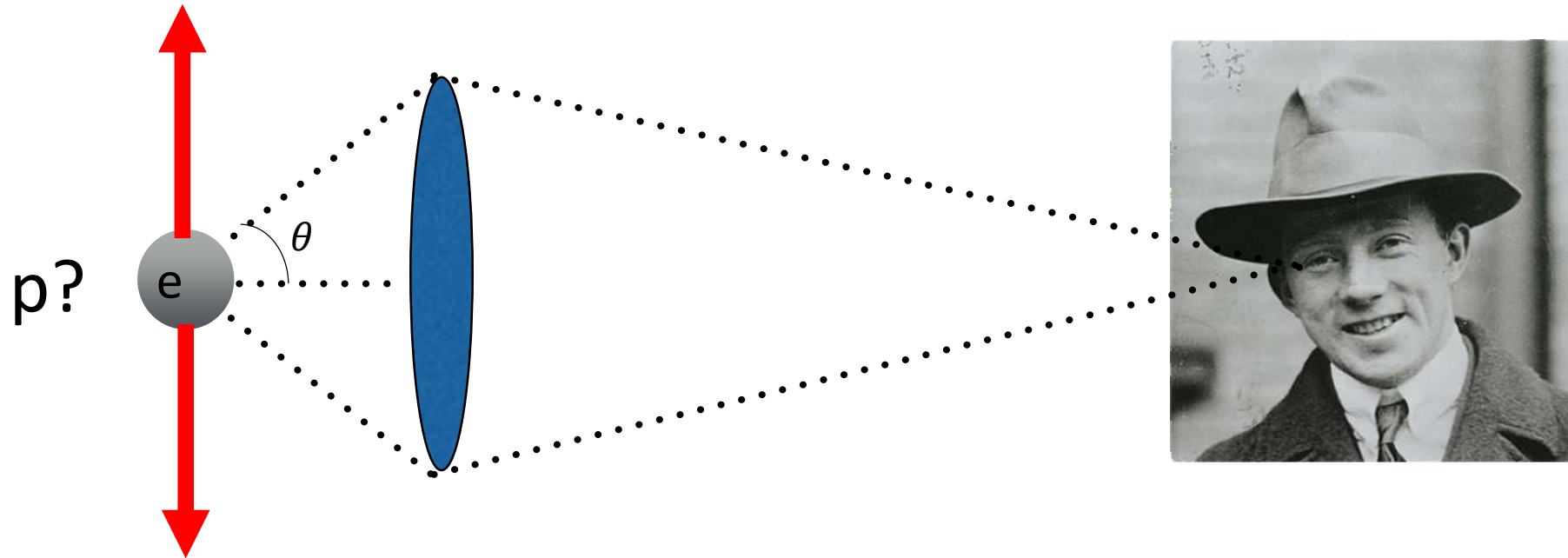
momentum kick
"Back-action"



Measurement back-action in Heisenberg's microscope

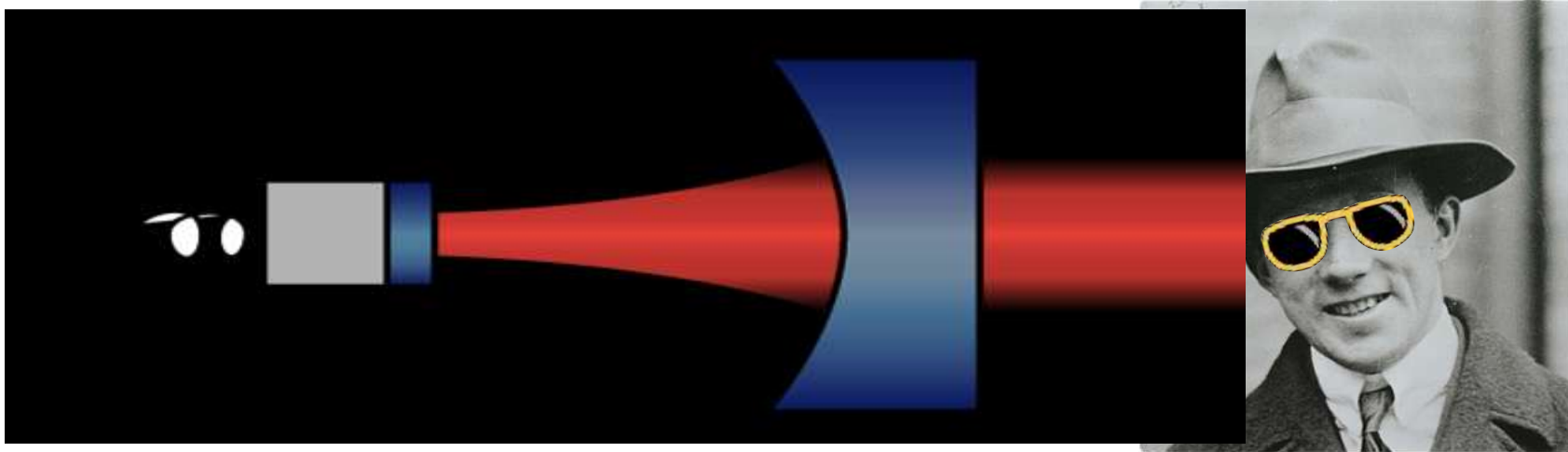


Measurement back-action in Heisenberg's microscope



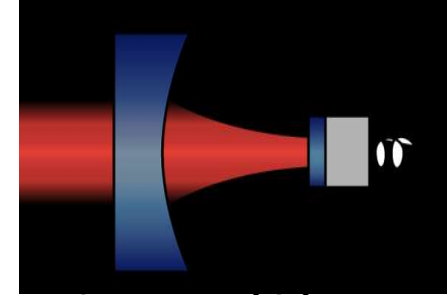
$$\Delta x \sim \frac{\lambda}{2 \sin \theta} \quad \& \quad \Delta p \sim \frac{h \sin \theta}{\lambda} \quad \rightarrow \quad \Delta x \cdot \Delta p \sim \frac{h}{2}$$

How about mechanical oscillator?



Standard quantum limit

an *ideal* measurement of mechanical motion

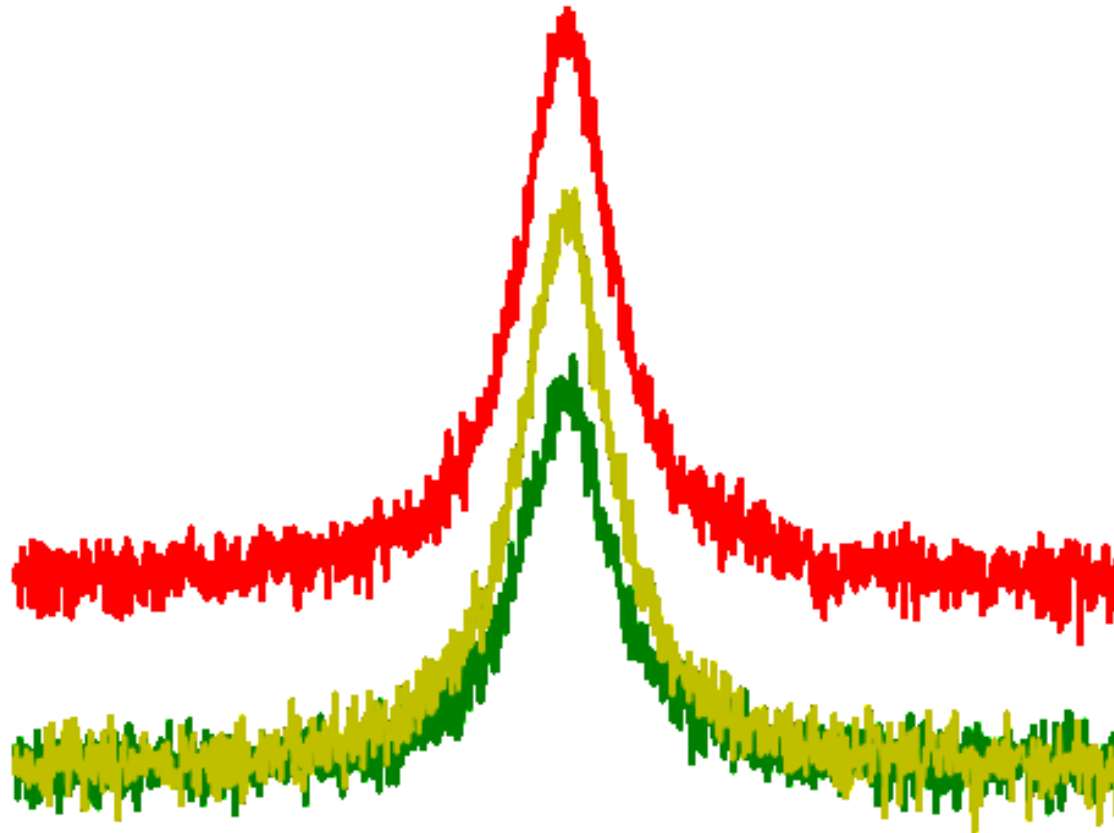


$$\hat{H} = \hbar\omega_c \hat{a}^\dagger \hat{a} + \hbar\omega_m \hat{b}^\dagger \hat{b} + \hbar g \hat{a}^\dagger \hat{a} (\hat{b}^\dagger + \hat{b})$$

photon mechanics or "phonon" interaction

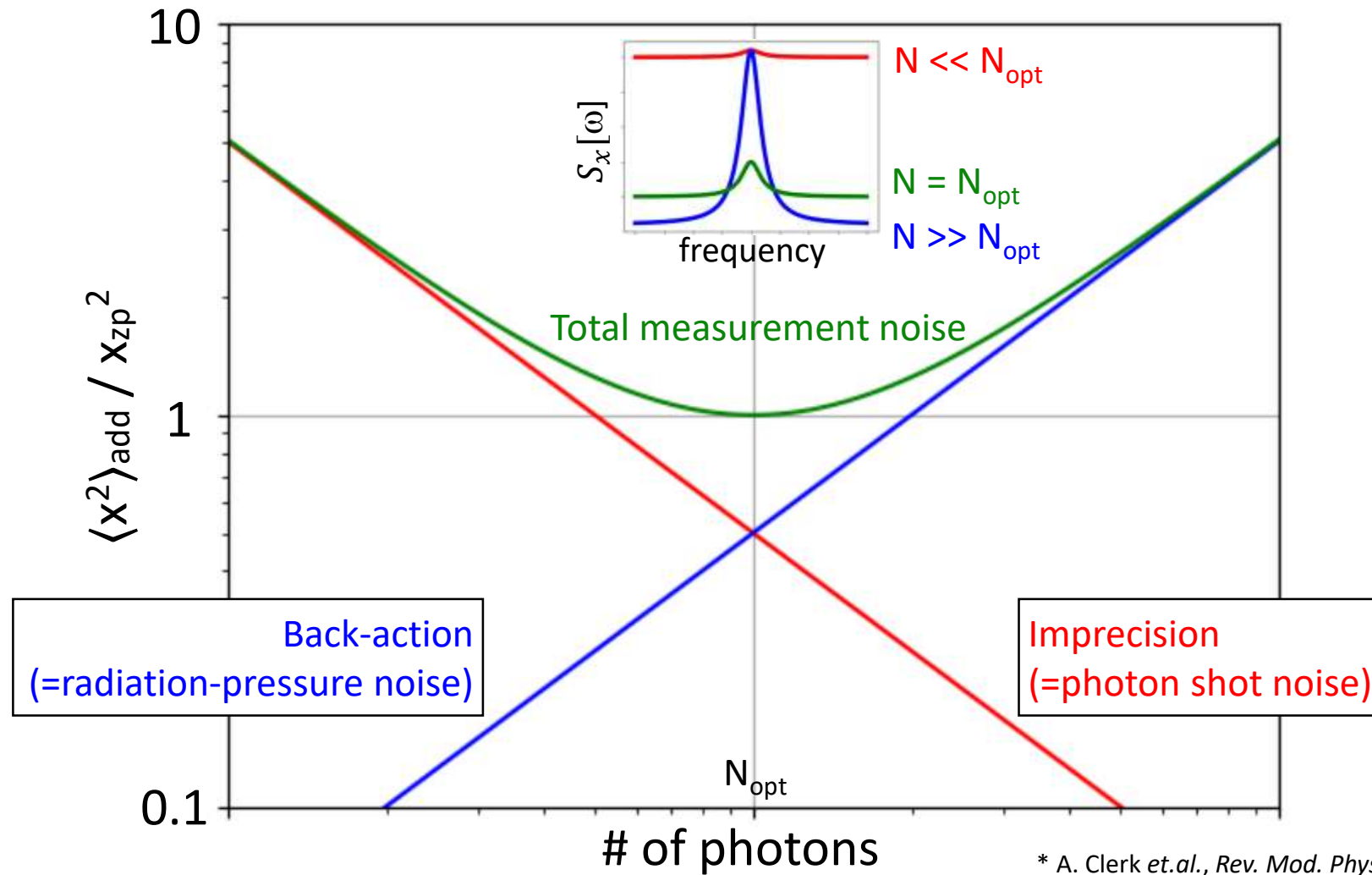
ation pressure noise
k-action"

on shot noise
recision"



* A. Clerk *et.al.*, *Rev. Mod. Phys.* **82**, 1155 (2010).

Standard quantum limit



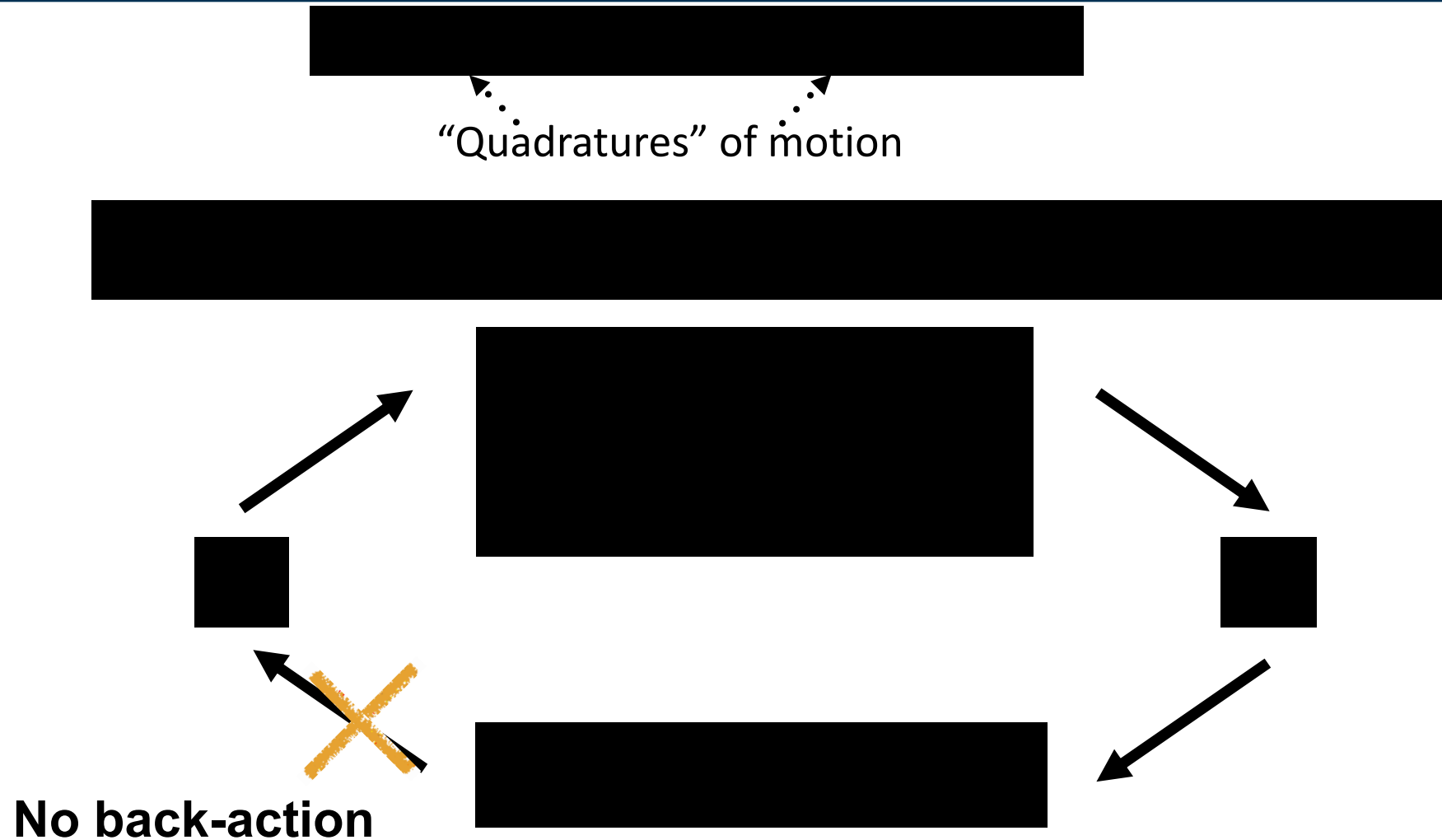
* A. Clerk *et.al.*, *Rev. Mod. Phys.* **82**, 1155 (2010).

Quantum limit in gravitational-wave detectors

Braginsky⁶ has pointed out that the above “quantum limits” on ΔX_1 , ΔX_2 , and ΔN pose serious obstacles for gravitational-wave detection: To encounter at least three supernovae per year, one must reach out to the Virgo cluster of galaxies. But gravitational waves from supernovae at that distance will produce $|\Delta X_1| \simeq |\Delta X_2| \lesssim 0.3 \times [m/(10 \text{ tons})] (\hbar/m\omega)^{1/2}$ in a mechanical oscillator on earth, corresponding to $\Delta N \lesssim 0.4(N + \frac{1}{2})^{1/2} [m/(10 \text{ tons})]$. For detectors of reasonable mass this signal is below the quantum limit.

* K. S. Thorne *et.al.*, *Phys. Rev. Lett.* **40**, 667 (1978).

Evading quantum back-action



Evading quantum back-action

$$\hat{x}(t) = \hat{X}_1 \cos \omega_m t + \hat{X}_2 \sin \omega_m t$$

“Quadratures” of motion

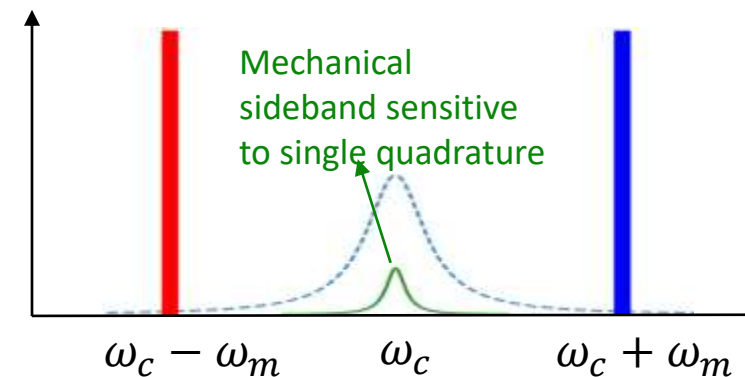
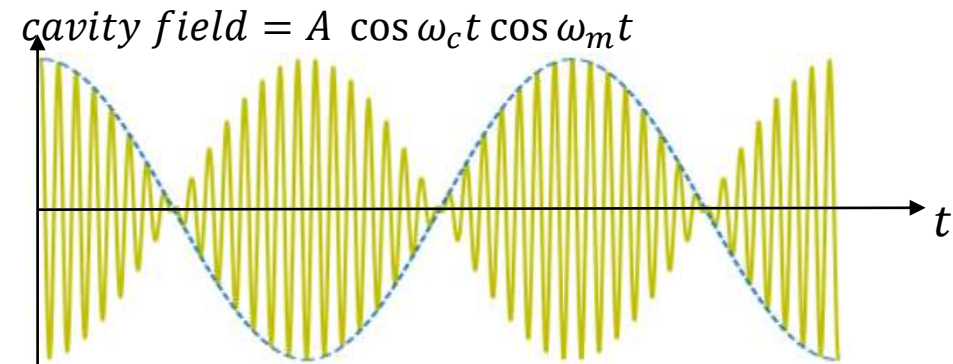
$\hat{X}_1, \hat{X}_2$: constants of motion of harmonic oscillator

⇒ can be measured with no back-action

⇒ back-action into the “unseen” quadrature

* Braginskii *et.al.*, *Sov. Phys. Usp.* **17**,644 (1975);

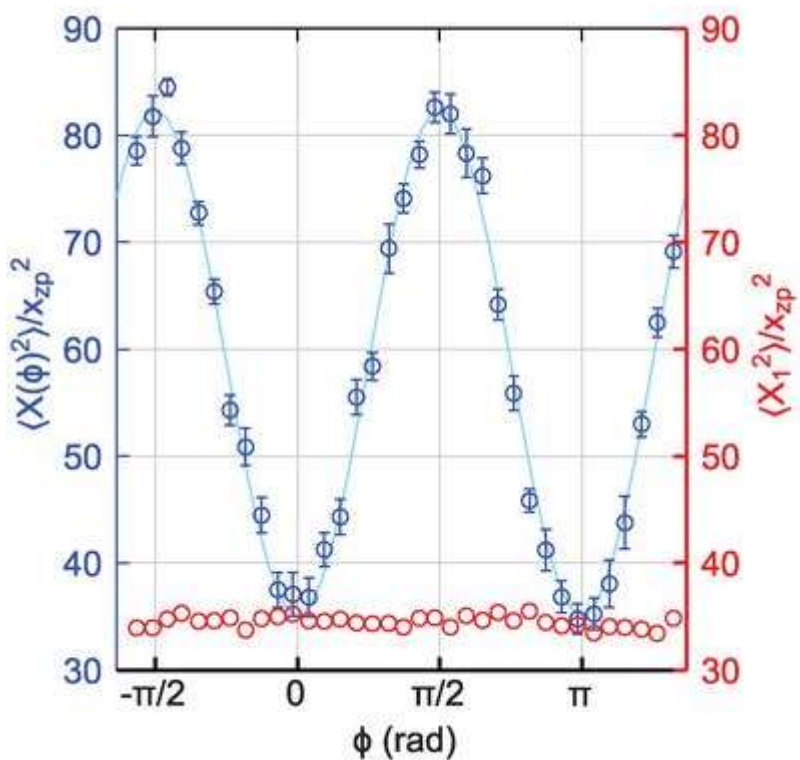
Thorne *et.al.*, *Phys. Rev. Lett.* **40**, 667 (1978).



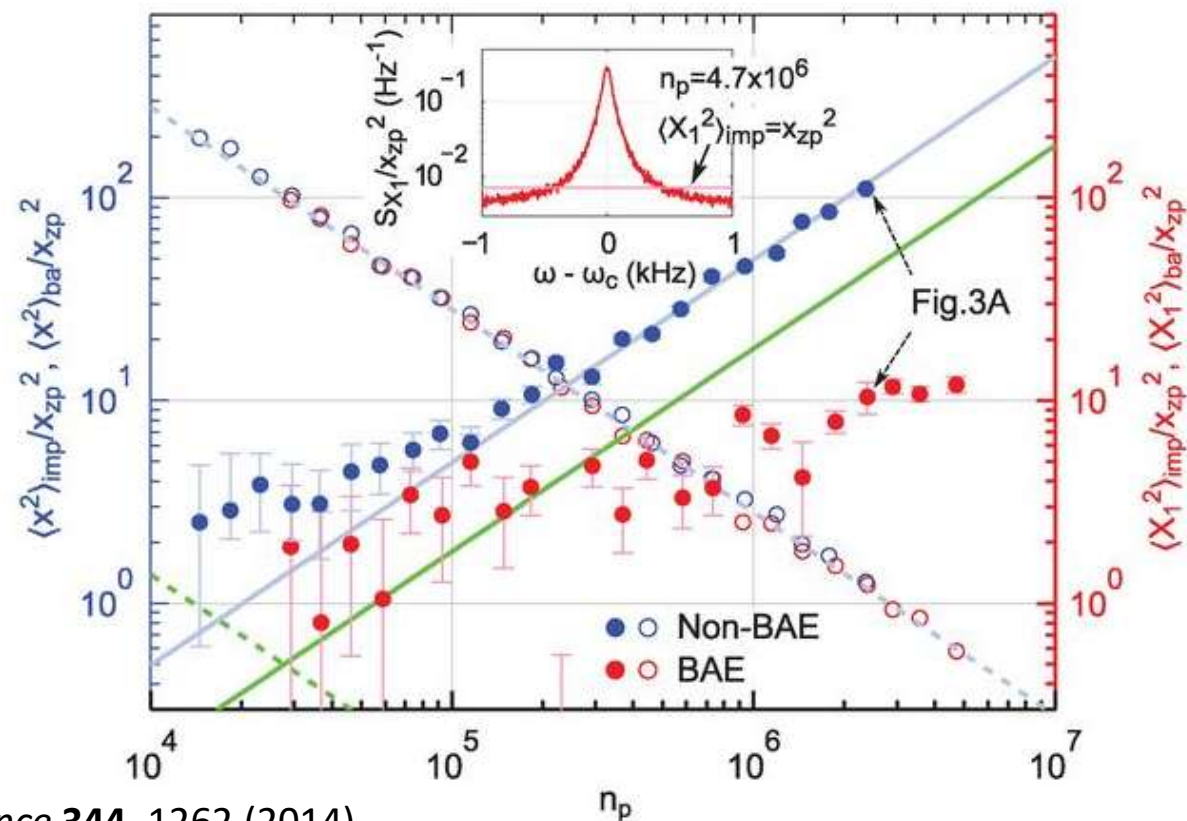
$$\hat{H}_{int} \propto \hat{X}_1 (1 + \cos 2\omega_m t) + \hat{X}_2 \sin 2\omega_m t$$

Experiments

back-action on ONE quadrature

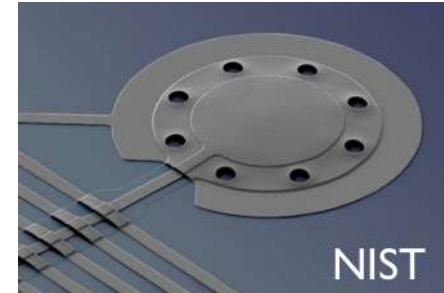
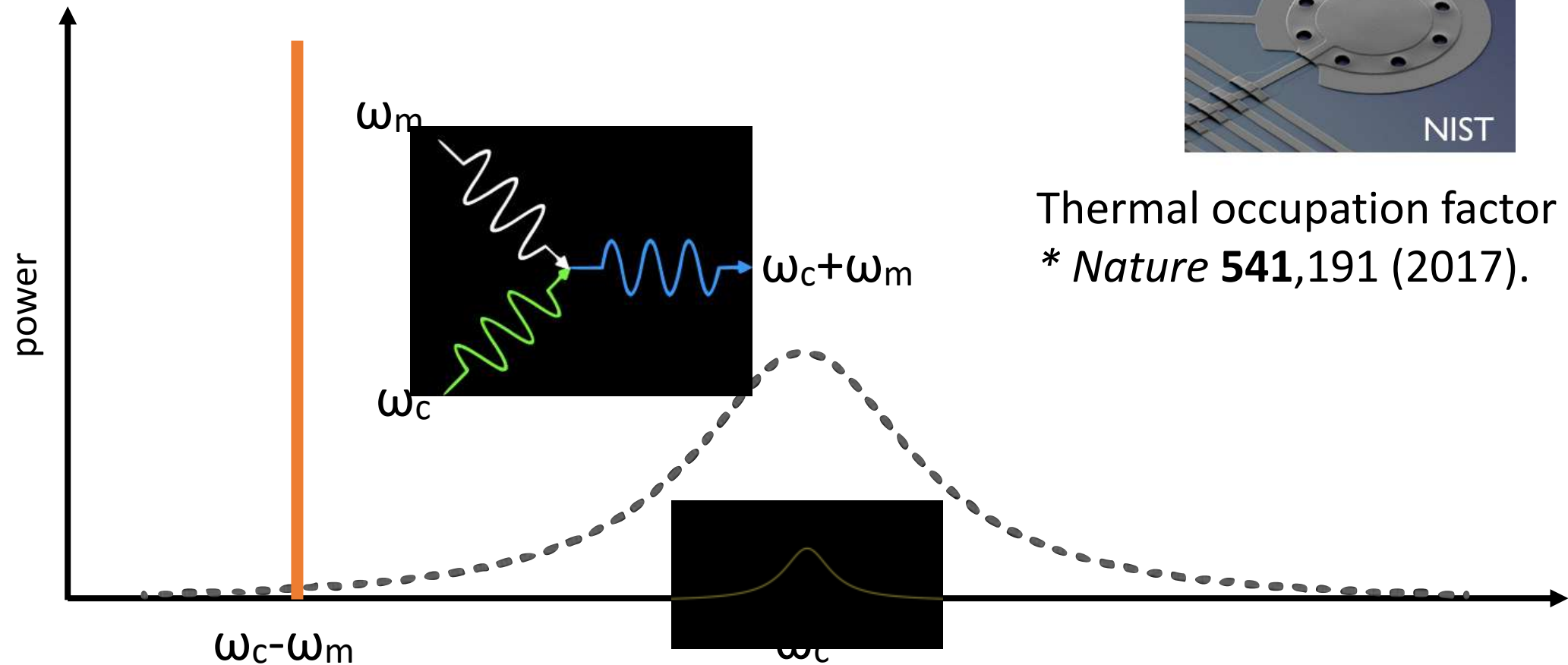


Evade quantum back-action by 8.5 dB



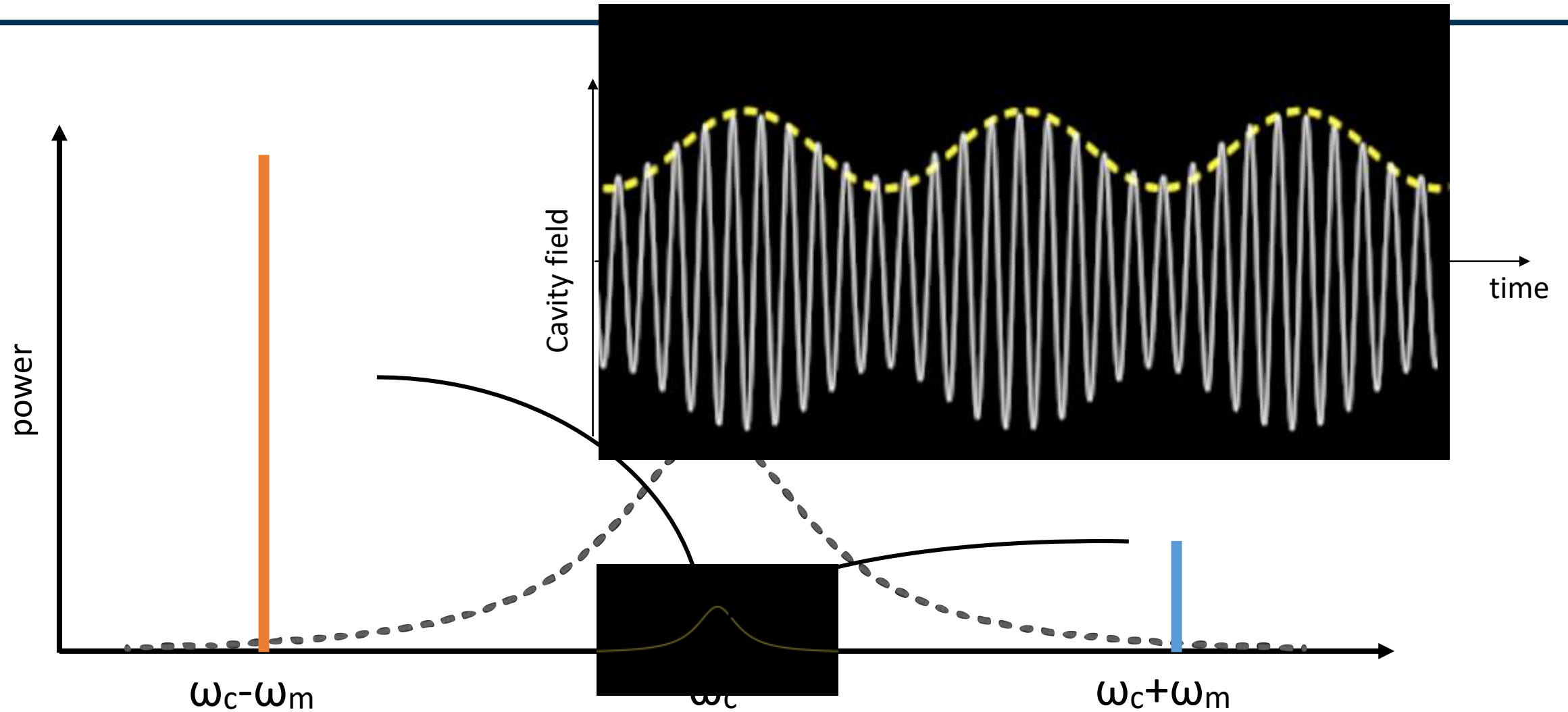
* JS et.al., *Science* **344**, 1262 (2014).

Ground state cooling of mechanical motion



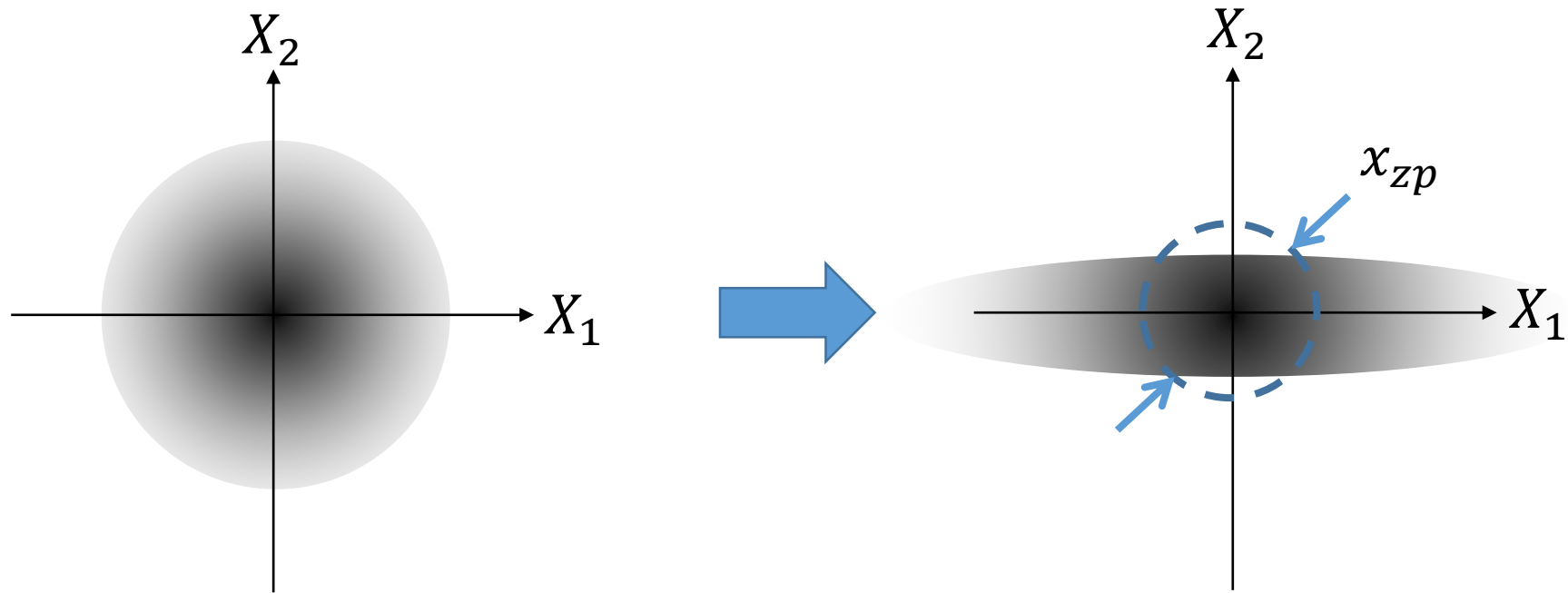
Thermal occupation factor ~ 0.2
* *Nature* **541**,191 (2017).

Phase-dependent cooling

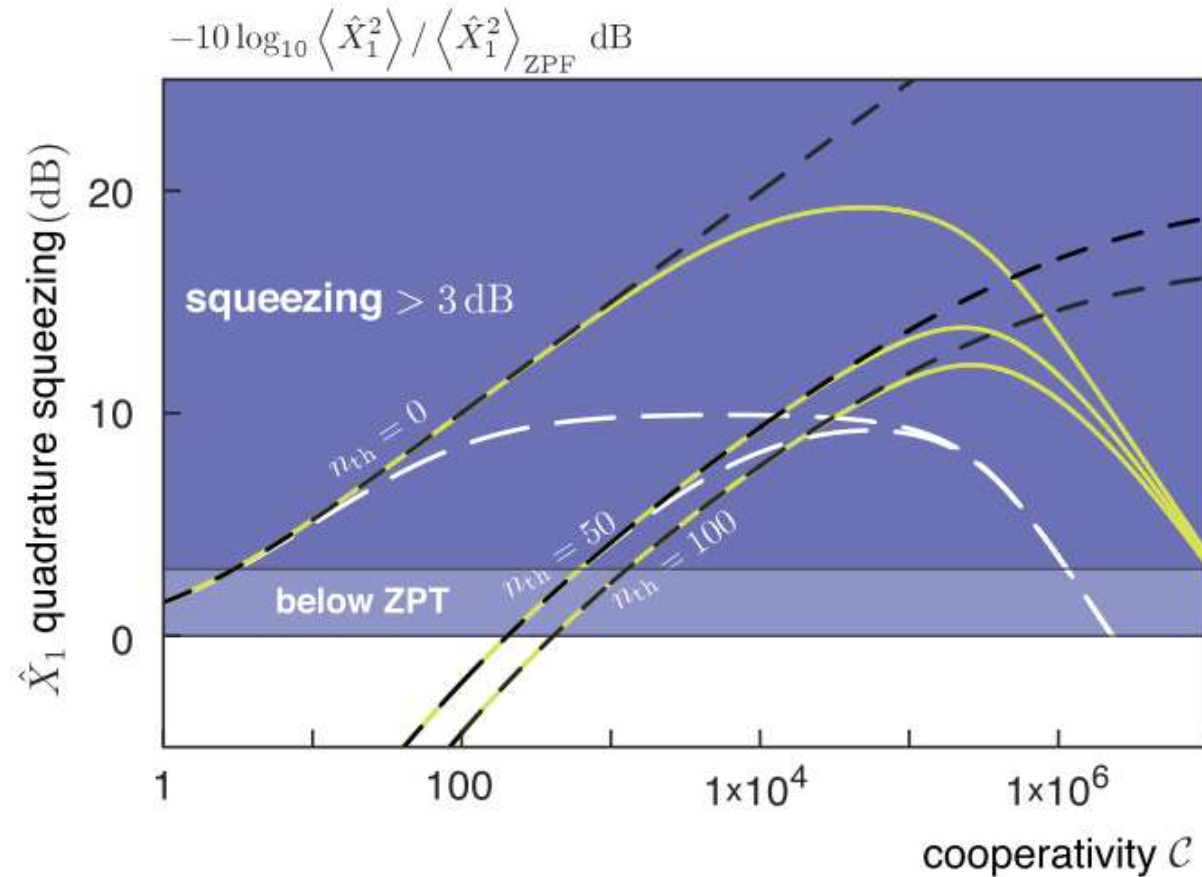


“Phase-dependent” reduction of mechanical motion (i.e. Squeezing)

$$\hat{x}(t) = \widehat{X}_1(t) \cos \omega_m t + \widehat{X}_2(t) \sin \omega_m t$$



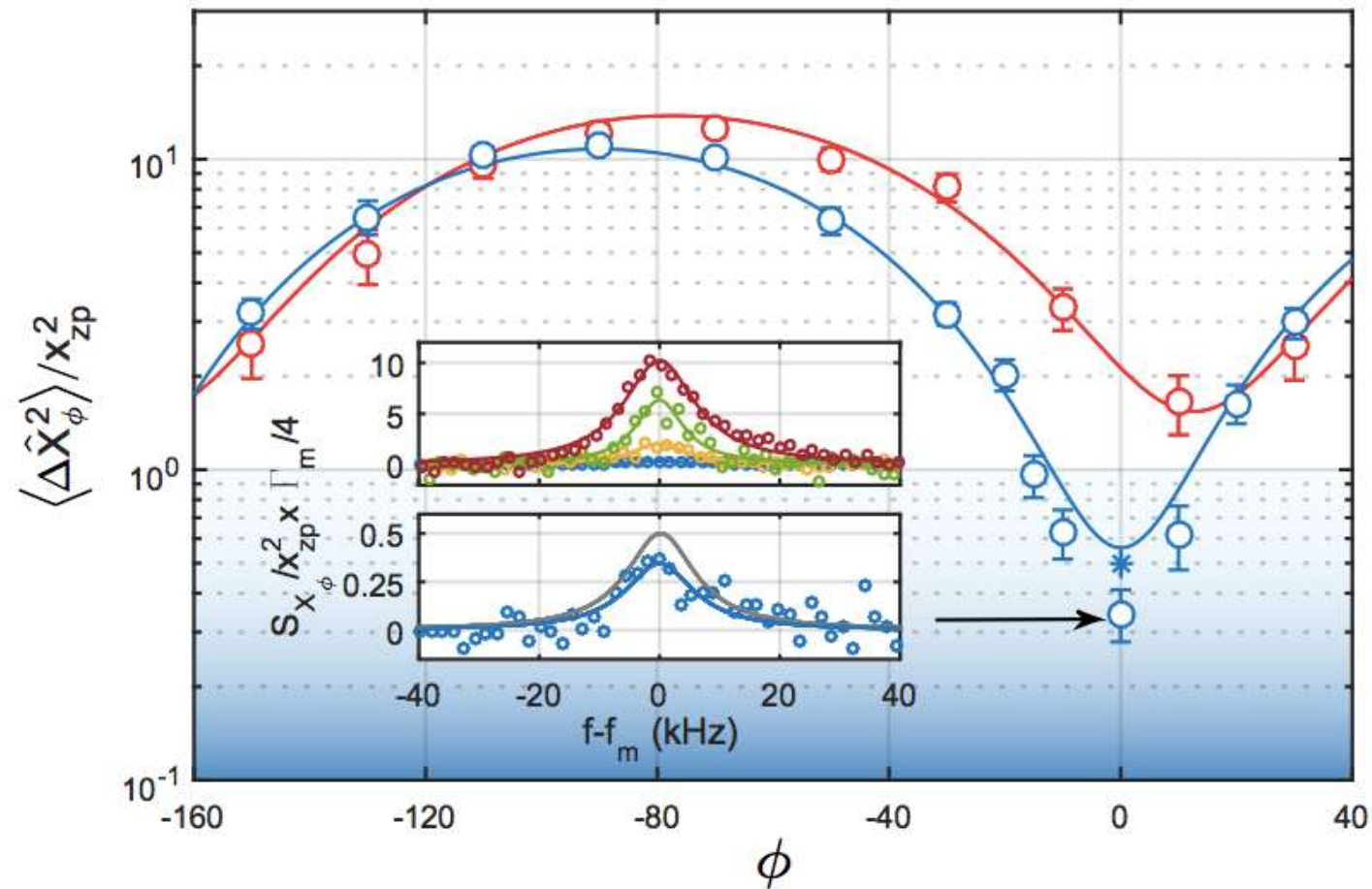
Arbitrarily large steady-state bosonic squeezing via dissipation



- Optimal ratio between red and blue power
- Squeezing beyond 3dB possible
- Steady state is squeezed thermal state
- State purity vs. squeezing

* Kronwald *et.al. Phys. Rev. A* **88**, 063833 (2014).

Squeezing more than 3 dB below zero-point

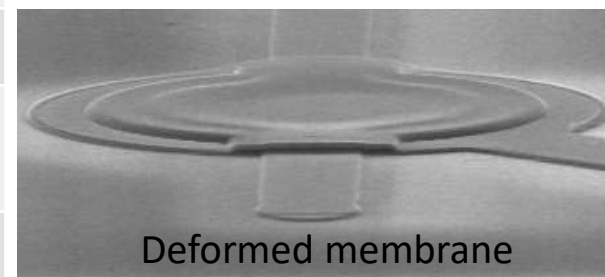


* Lei, Weinstein, **JS**, Wollman, Kronwald, Marquardt, Clerk, Schwab, *PRL* **117**, 100801 (2016).

Niobium for cavity optomechanical sensing under magnetic field

Niobium superconducts at higher temperatures and magnetic fields.

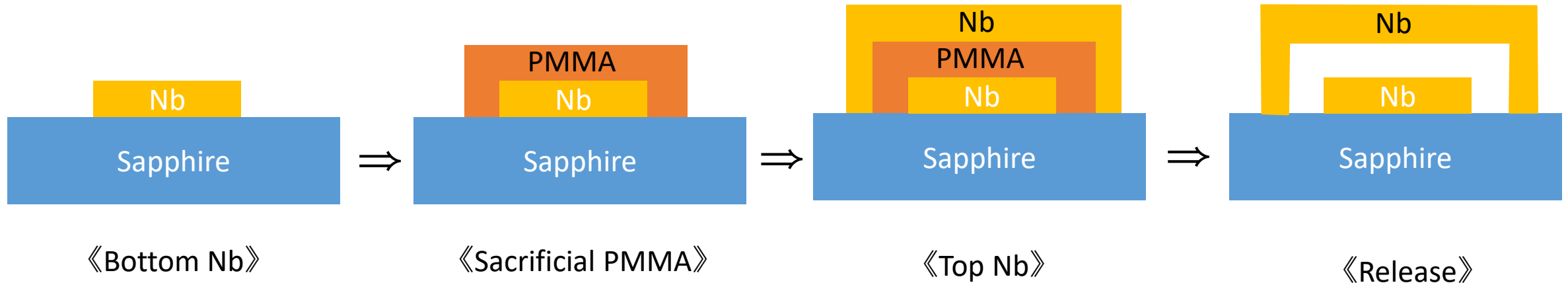
	Aluminum	Niobium
Critical Temperature (T _c)	1.2K	9.26K
Critical Magnetic Field(H _c)	0.01 T	0.82 T
Density	2700 kg/m ³	8570 kg/m ³
Young's modulus	70 Gpa	105 GPa
Poisson ratio	0.35	0.4
Advantages	• Easy to control the film stress • Large zero point motion due to the small mass	• Good mechanical properties • High critical temperature and magnetic field
Disadvantages	• Low critical temperature	• Difficult to control the film stress



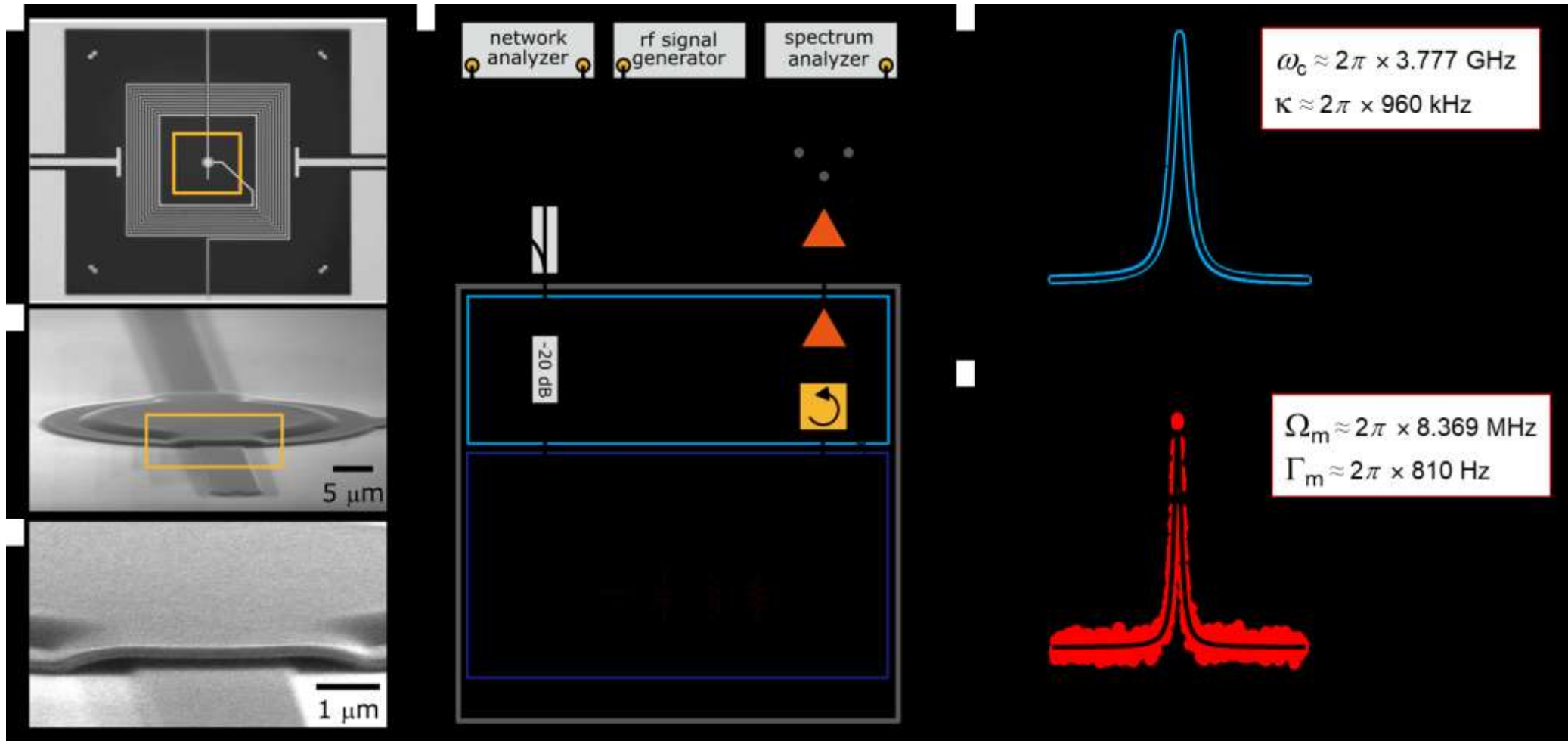
Jinwoong Cha (KRISS)

* J. Cha *et.al.*, “Superconducting Nanoelectromechanical Transducer Resilient to Magnetic Fields”, *Nano Letters* **21**, 1800 (2021).

Fabrication

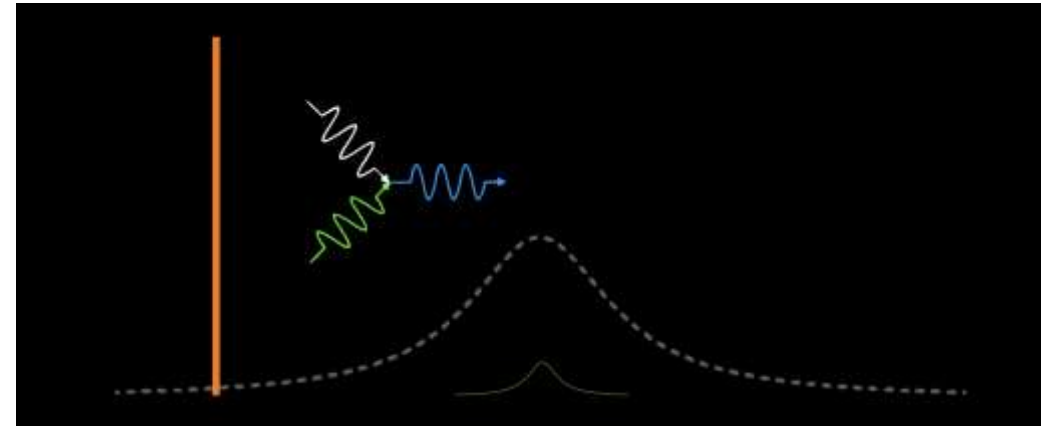
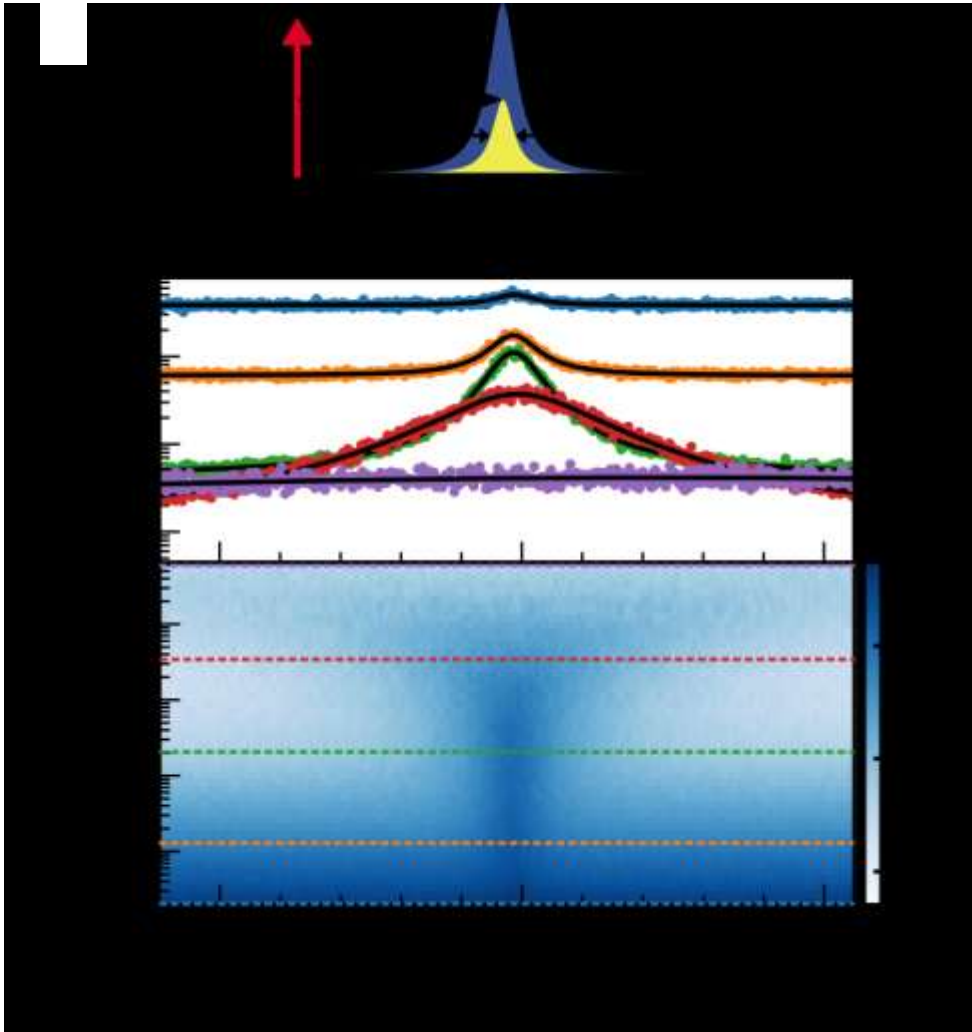


Microwave cavity optomechanics at 4.2 K



* Cha *et.al*, *Nano Letters* **21**, 1800 (2021).

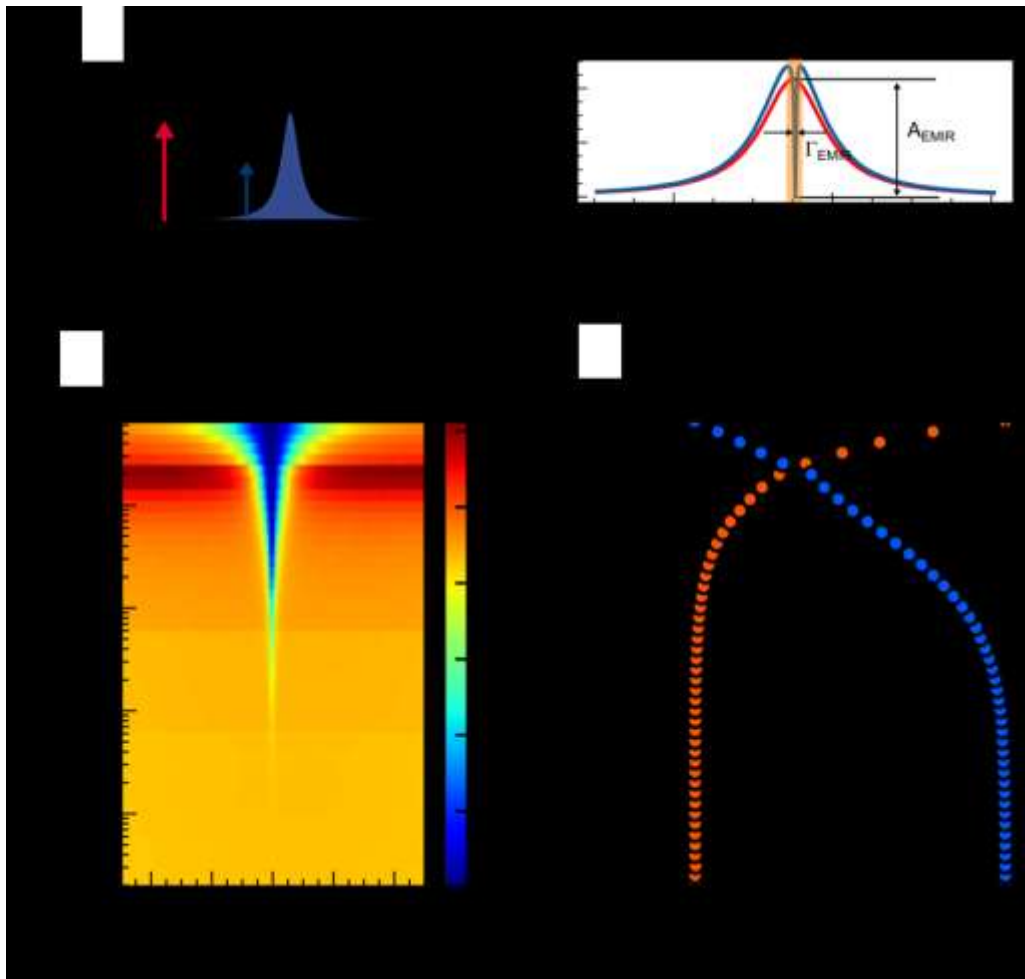
Back-action cooling at 4.2 K



- Cooling process accompanies with mechanical linewidth broadening
- Efficient cooling of mechanical mode temperature from 4.2 K to 76 mK

Cha *et al.*, *Nano Letters* **21**, 1800 (2021).

Electromechanical induced reflection of microwave



- Probe microwave interferes destructively with mechanical sideband from pump

- Reflection window

$$\Gamma_{EMIR} = \Gamma_m \left(1 + \frac{4g_0^2 n_d}{\kappa \Gamma_m} \right) = \Gamma_m (1 + C)$$

- Single photon coupling

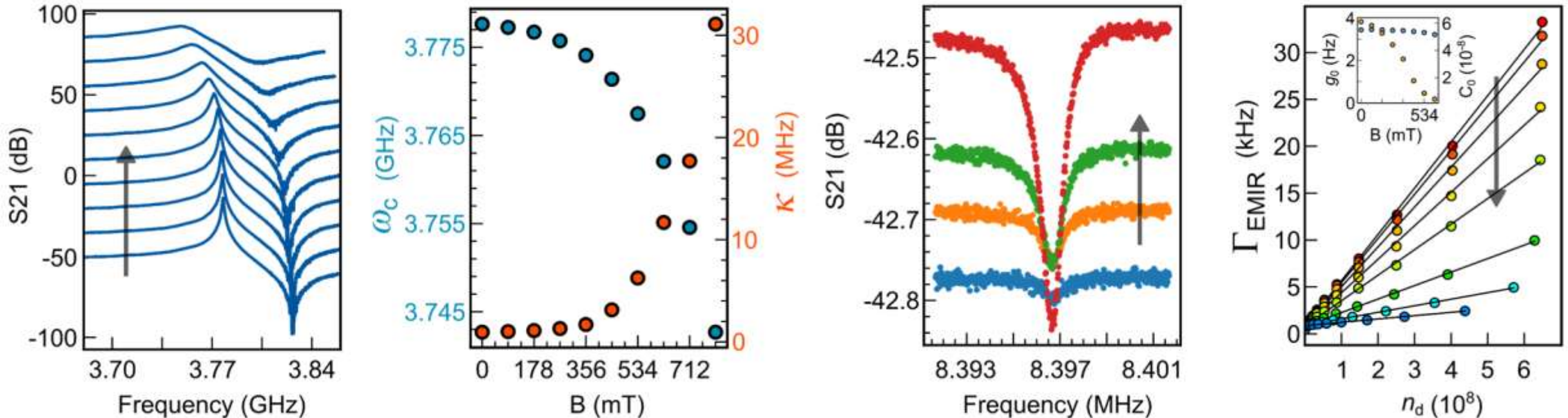
$$g_0 \approx 3.3 \text{ Hz}$$

- Cooperativity

$$C \approx 40$$

* Cha *et.al*, *Nano Letters* **21**, 1800 (2021).

Operation in magnetic field



- Magnetic field B affects the microwave resonance frequency and linewidth.
- Mechanical sideband signal persists even at 0.8 T.
- Cooperativity decreases as B increases due to the increasing cavity decay rate.
- Single-photon coupling rate is independent of magnetic field.

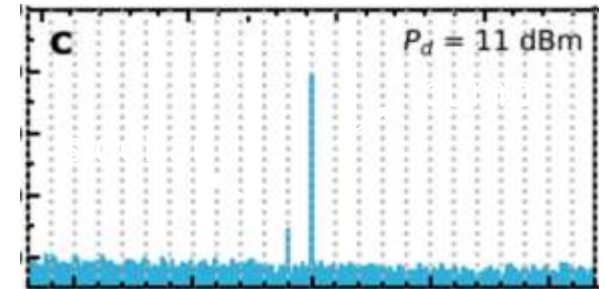
* Cha *et.al*, *Nano Letters* **21**, 1800 (2021).

Niobium optomechanics for non-linear optomechanics

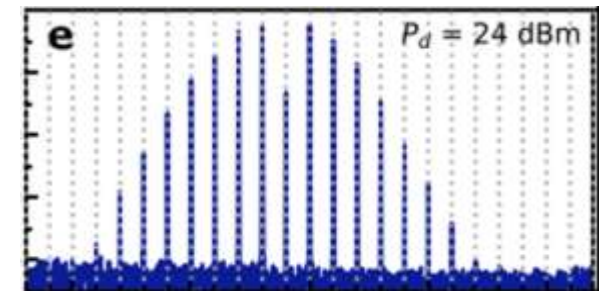
Nb handles more RF current ($\sim 6 \times 10^3$ more RF photons)

	Aluminum	Niobium
Critical Temperature (Tc)	1.2K	9.26K
Critical Magnetic Field(Hc)	0.01 T	0.82 T
Density	2700 kg/m ³	8570 kg/m ³
Young's modulus	70 Gpa	105 GPa
Poisson ratio	0.35	0.4
Advantages	• Easy to control the film stress • Large zero point motion due to the small mass	• Good mechanical properties • High critical temperature and magnetic field
Disadvantages	• Low critical temperature	• Difficult to control the film stress

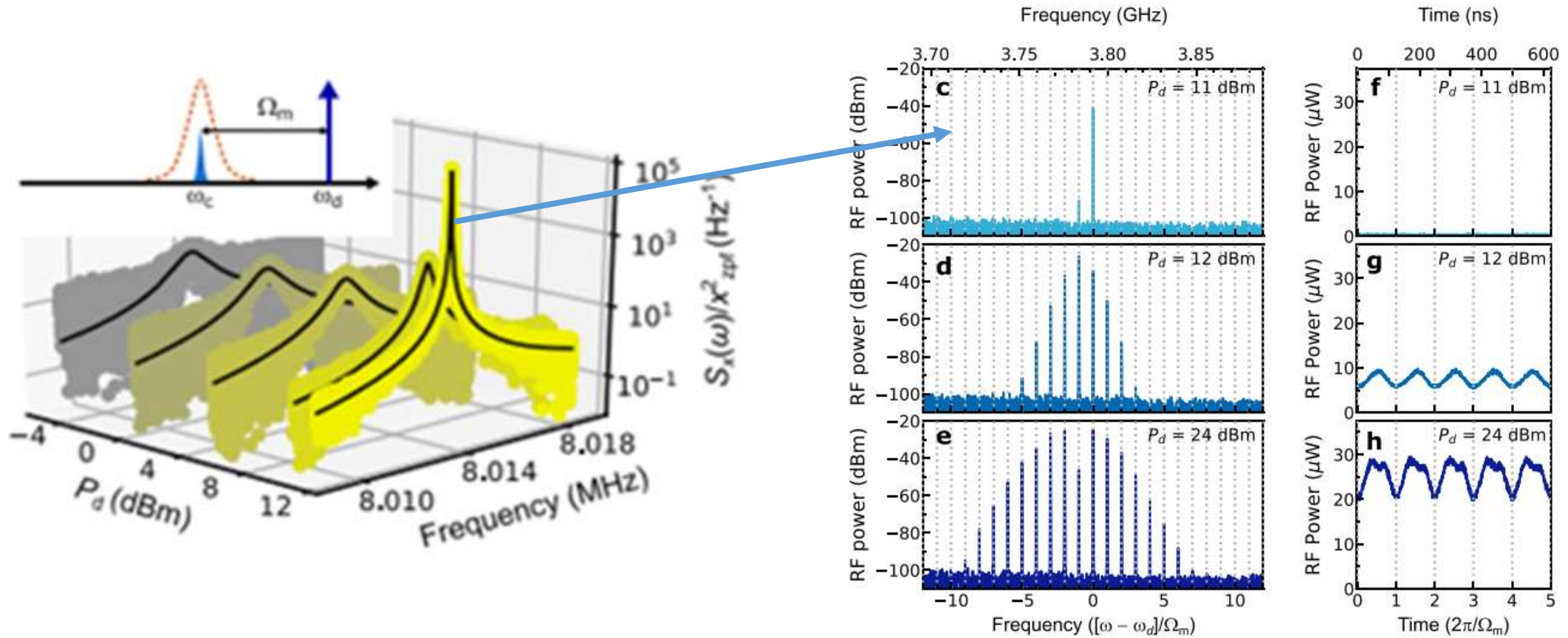
Linear photon-phonon conversion



Multiple non-linear sidebands



Nonlinear responses above instability



* J. Shin *et.al.*, *Nano Letters* **22**, 5459 (2022).

Optomechanical frequency comb

New J. Phys. **20** (2018) 043013

<https://doi.org/10.1088/1367-2630/aab5c6>

Optomechanical frequency combs

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¹ Department of Electrical and Computer Engineering, The University of Texas at Austin, Austin, TX 78712, United States of America

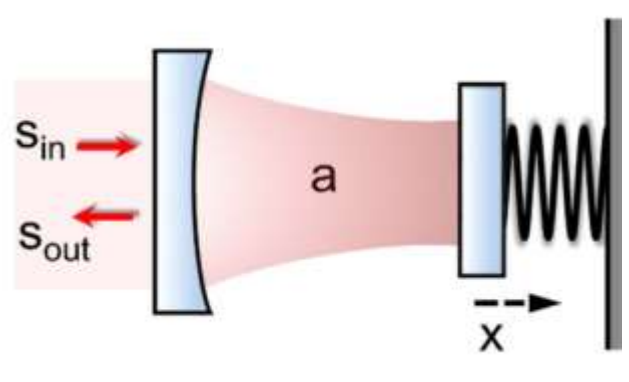
² Photonics Initiative, Advanced Science Research Center, City University of New York, New York 10031, United States of America

³ Physics Program, Graduate Center, City University of New York, New York 10016, United States of America

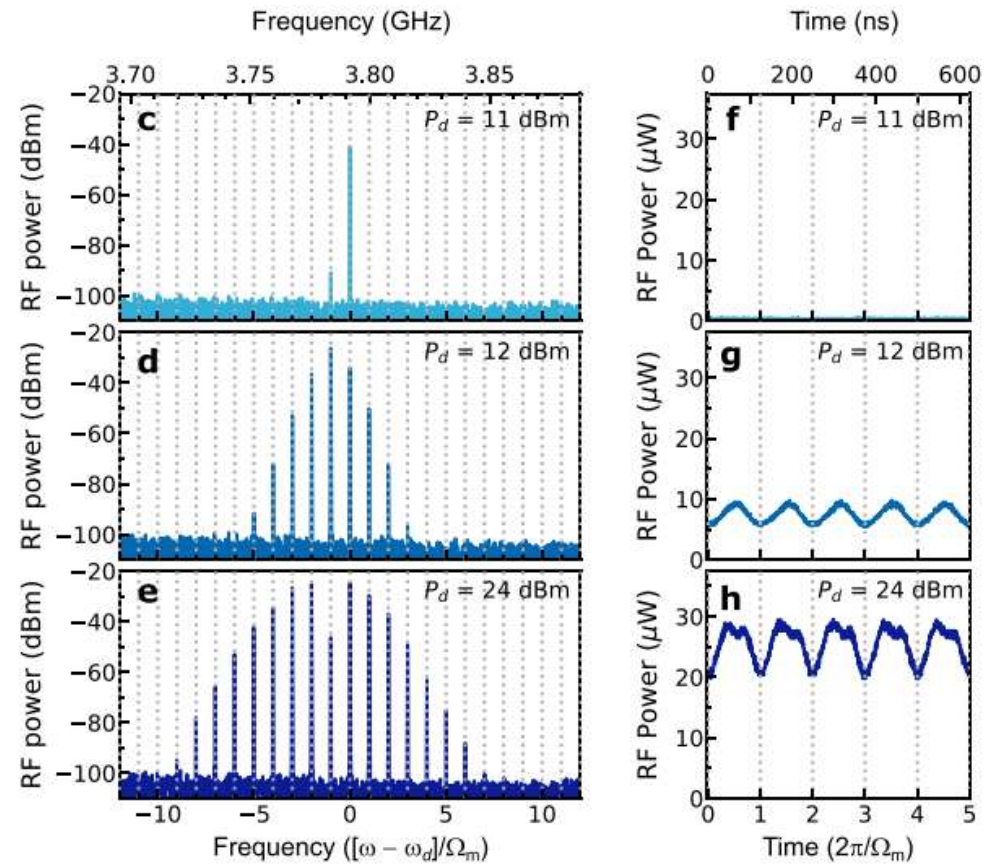
⁴ Department of Electrical Engineering, City College of The City University of New York, New York 10031, United States of America

E-mail: aalu@gc.cuny.edu

Keywords: combs, optomechanics, nanophotonics

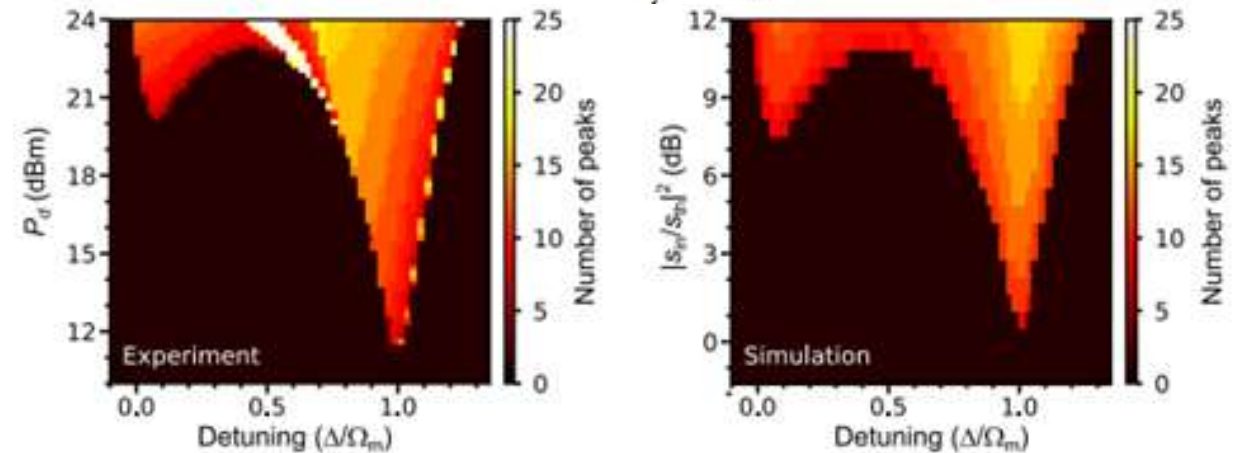

$$\frac{da}{dt} = \left(i(\Delta + Gx) - \frac{\kappa}{2} \right) a + \sqrt{\kappa_e} s_{\text{in}}$$
$$\frac{d^2x}{dt^2} + \Gamma_m \frac{dx}{dt} + \Omega_m^2 x = \frac{\hbar G}{m} |a|^2$$

Optomechanical frequency comb

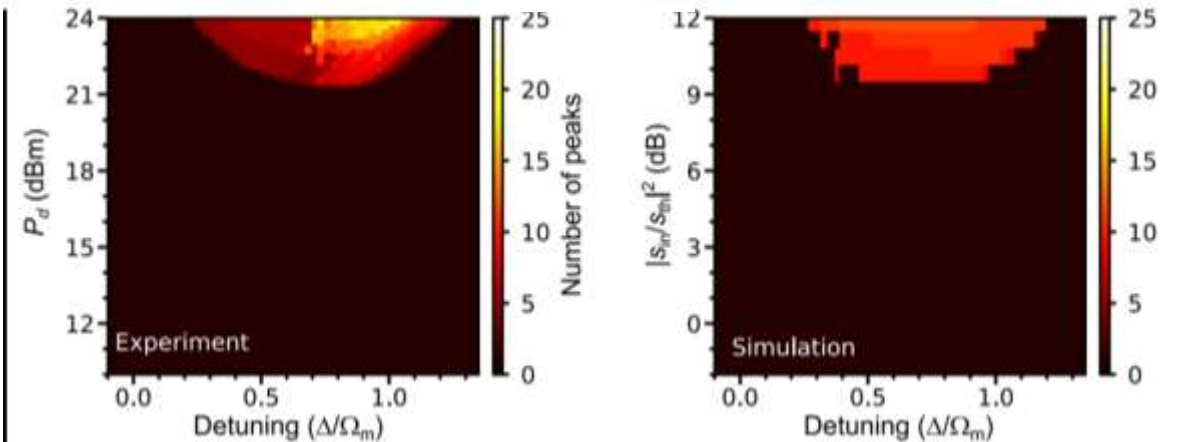


* J. Shin *et.al.*, *Nano Letters* **22**, 5459 (2022).

$$\kappa \sim 0.14\Omega_m$$

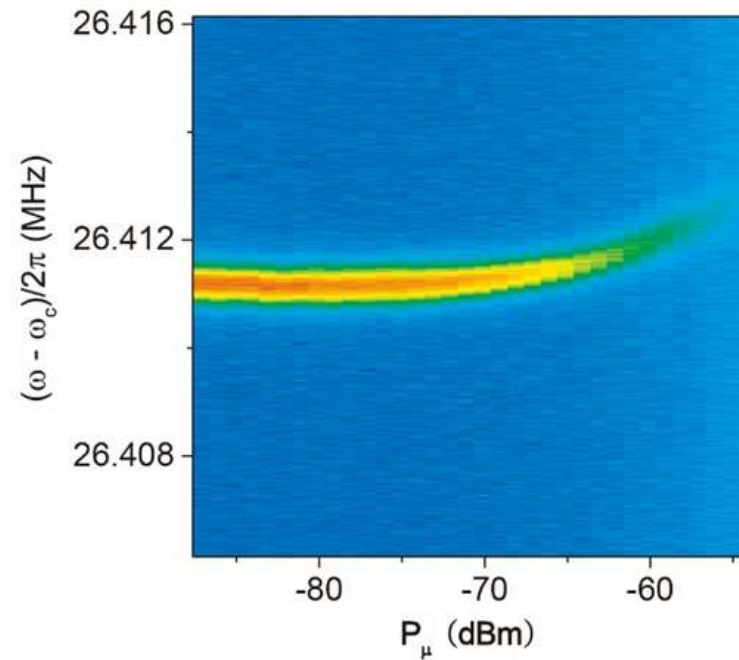
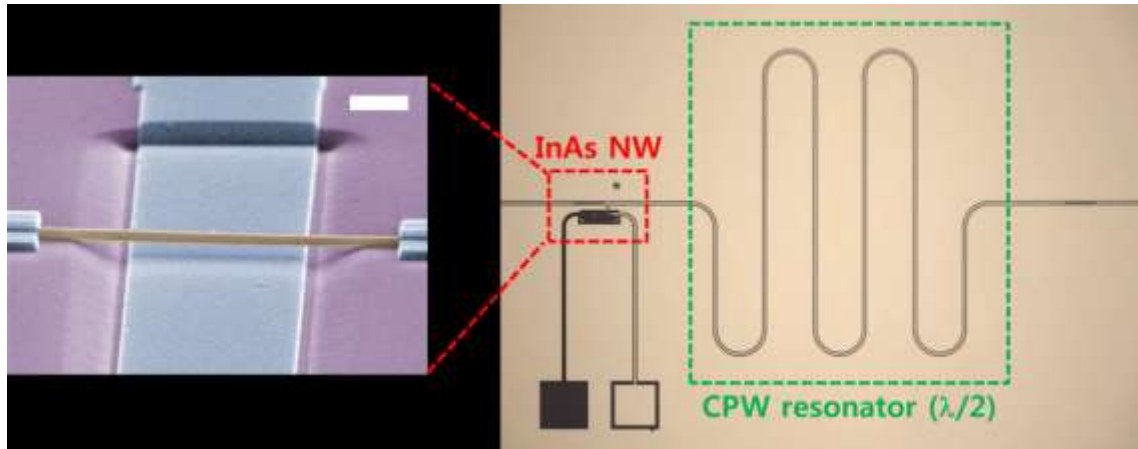


$$\kappa \sim \Omega_m$$



Utilizing Dissipation: Nanomechanical microwave bolometer

Nanomechanical sensor detects heat from microwave photons

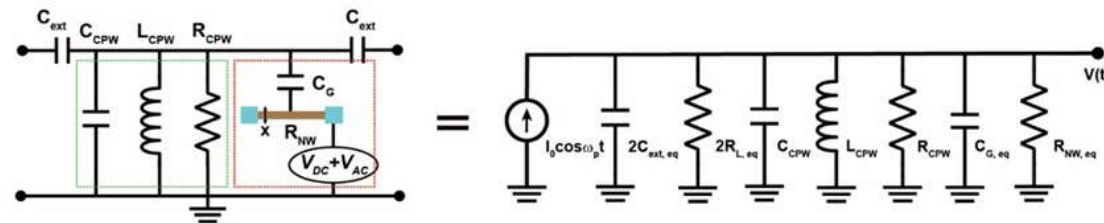
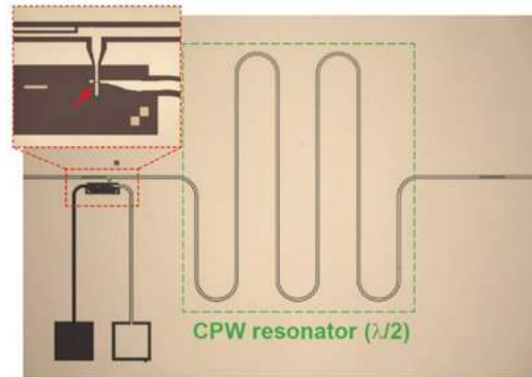
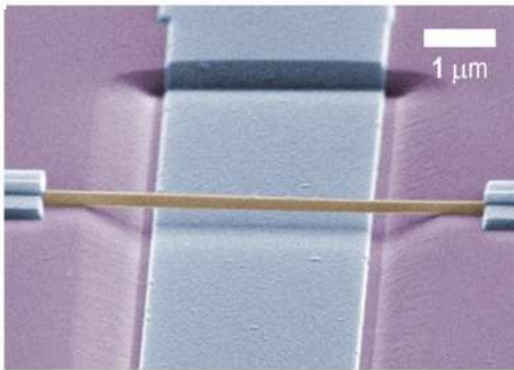


Jihwan Kim
(now at ADD)

* J. Kim *et.al.*, “Nanomechanical Microwave Bolometry with Semiconducting Nanowires”, *Physical Review Applied* **15**, 034075 (2021).

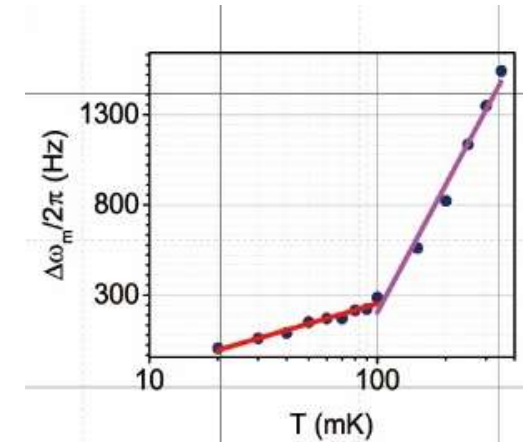
InAs nanowire based cavity optomechanics

Resistive nanowire dissipates microwave power

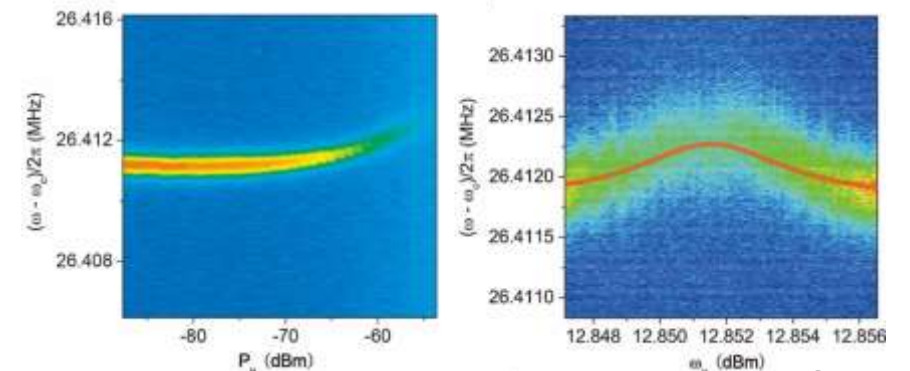


* J. Kim *et.al.*, *Physical Review Applied* **15**, 034075 (2021).

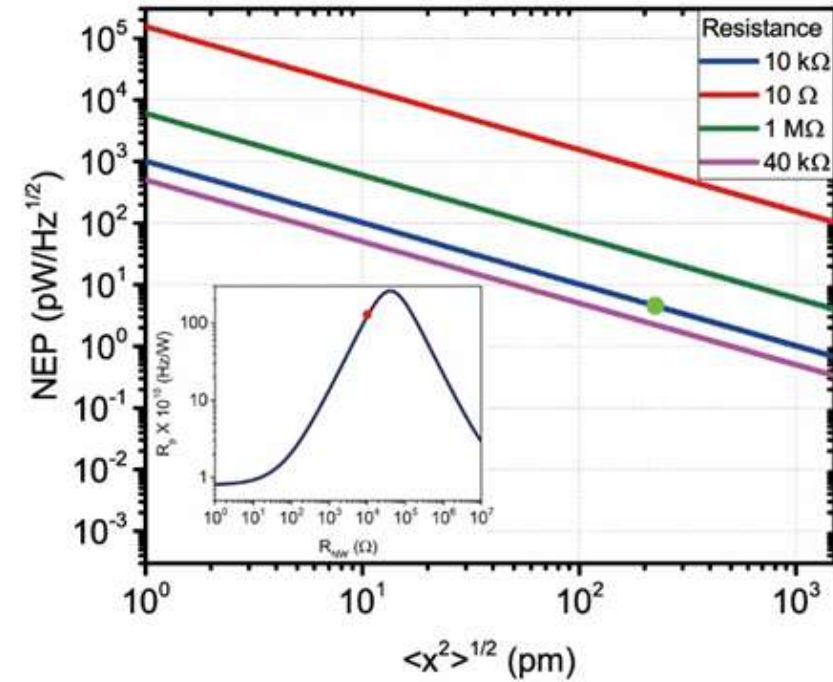
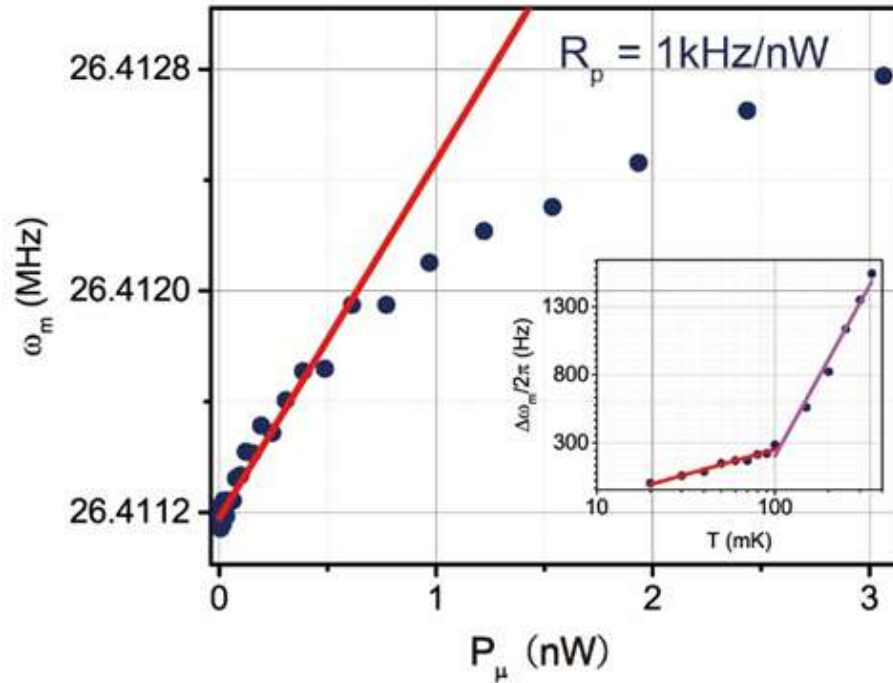
Nanomechanical thermometer



Mechanical resonance senses microwave power



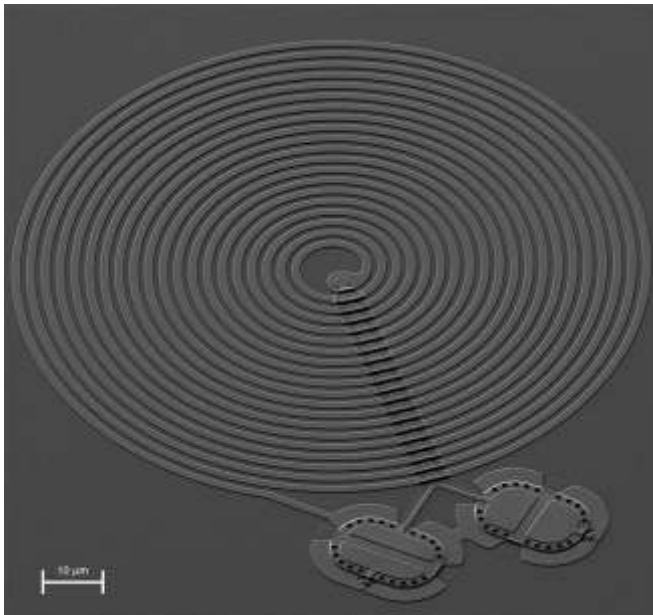
Nanomechanical microwave bolometry



- “Noise equivalent power” $\text{NEP} = 4.5 \text{ pW/Hz}^{1/2}$
- Maximum detectable power $\sim \text{nW}$
- c.f. Josephson bolometer has $\text{NEP} \sim \text{aW/Hz}^{1/2}$ and maximum power $\sim \text{fW}$ (ref. *Nature* **586**, 42 (2020))

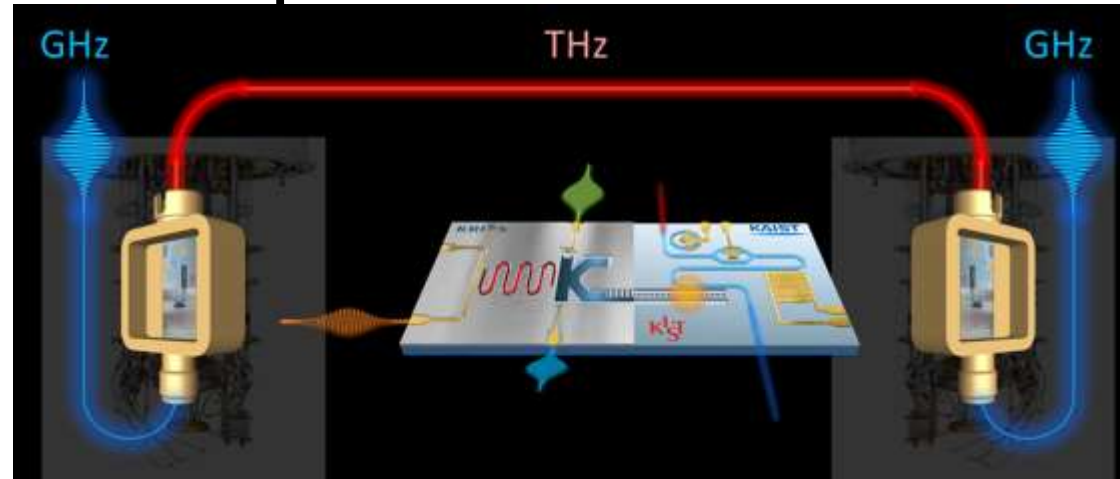
Outlook: quantum transducer and sensors

entangled force sensors



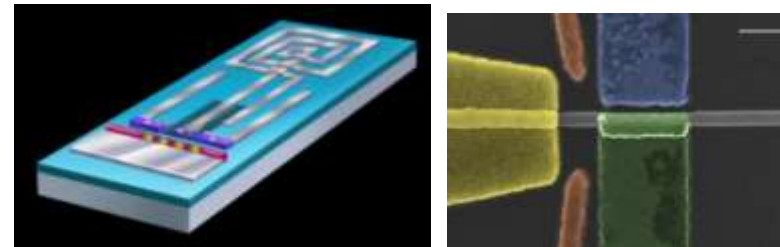
*Kotler *et al.*, *Science* **372**, 622 (2021).

quantum transduction



* “Integrated quantum interconnects for long-distance quantum networks” funded by NST

sensors for new physics



* “Quantum electromechanical interface for Majorana qubits” funded by Samsung foundation

중시계 여름학교

Hybrid Quantum Device Lab

hql.postech.ac.kr

